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### **ORIGINAL ARTICLE**

# Deuterium-bearing molecules in cold and warm dense interstellar clouds

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**Abstract** We have studied the production of key observed deuterium species for TMC-1, Orion, L134N and W3IRS4 clouds, of densities  $10^4$ ,  $10^4$ ,  $5 \times 10^4$ , and  $10^6$  cm<sup>-3</sup>, respectively, by using the pseudo-time dependent gas-phase chemistry. These clouds have temperatures in the range 10–70 K. The main results by using the more extensive chemical network with the most updating reaction rates show that the most of calculated fractional abundances are in agreement with observations, and suggest that triply-deuterated ammonia could be detectable in dark clouds. Also our models show that large abundance of NH<sub>2</sub>D and NHD<sub>2</sub> can be produced in the interiors of cold dense clouds at steady state time.

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#### 1. Introduction

Because deuterium-bearing molecules used to: probes of the physics of interstellar clouds, study the relation of connection between interstellar and cometary ices and understand the formation mechanism of isotopic composition of interstellar molecules (Roberts and Millar, 2000a; Robets et al., 2002; Shah and Wotten, 2001; van der Tak et al., 2002). Therefore many theoretical and observational studies which have concerned on the formation of deuterated molecules in interstellar clouds.

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In the last few years the studying of doubly-deuterated molecules in the interstellar medium has gained considerable attention. This is due to a large amount of doubly-deuterated formaldehyde, D<sub>2</sub>CO, and ammonia, NHD<sub>2</sub>, have been observed towards low mass protostar IRAS16293-2422 (Ceccarelli, 1998), Orion cloud (Turner, 1990), very young protostellar core 16293E (Loinard et al., 2001) and in the molecular cloud L1689N (Ceccarelli et al., 2002; Loinard et al., 2001). Also NHD<sub>2</sub> observed in the molecular cloud L134N (Roueff et al., 2000). While triply-deuterated ammonia, ND<sub>3</sub> observed in the low mass protostar NGC1333-IRAS4 (van der Tak et al., 2002), in the dark cloud B1 (Lis et al., 2002). But the methanol observed in the low mass protostar IRAS16293-2422 (Parise et al., 2002). All these observations suggest that deuteration of formaldehyde and methanol is produced during the cold and dense cloud (Ceccarelli et al., 2001).

There are two chemical networks for the formation of the deuterated molecules:

First, some deuterium-bearing molecules can be formed by the gas-phase reactions (Millar et al., 1989; Rodgers and Millar, 1996; Roberts and Millar, 2000a; Robets et al., 2002). At low temperatures 10–70 K, the  $H_2D^+$  species is the key of gas-phase reaction to form DCO<sup>+</sup> and  $N_2D^+$  (van der Tak et al., 2002). By gas-phase network, Lis et al. (2002) explained the observed abundance ratio of ND<sub>3</sub> in B1 cloud. Similarly, Tine et al. (2000) investigated the fractional abundance of NH<sub>2</sub>D, which observed in L183 and TMC-1 clouds.

Second, the surface chemistry, in which deuterated molecules can be formed on dust grains. The high abundances of HDCO,  $D_2CO$  and  $CH_3OD$  in warm clouds are derived from the occurrence of grain surface chemistry during an earlier cold era followed by evaporation into the gas as temperature rise (Ceccarelli et al., 2001; Caselli et al., 2002; van der Tak et al., 2002; Parise et al., 2002). The formation of singly- and doubly-deuterated isotopomers of formaldehyde and singly-, doubly- and multiply-deuterated isotopomers of methanol on dust grain has been studied by Caselli et al. (2002), with a semi-empirical modified rate approach and Monte Carlo method in temperature range 10–20 K.

In the present work, we shall confine ourselves to calculate the chemical abundances of key observed deuterium species, in several different interstellar clouds and comparison with both available observation and other theoretical models which are given for justification. This paper is organized as follows, in Section 2, chemical models are given. Section 3 contains a brief description of our gas phase chemical models. The conclusions are given in Section 4.

#### 2. Chemical models

In order to interpret the behavior of deuterium-bearing molecules on cold and warm interstellar clouds, we have carried a number of a pseudo-time dependent chemical models, which calculate the varying abundances of 408 species (130 of them containing deuterium) linked by 5320 reactions.

Our gas-phase model considers a standard gas-phase chemistry, in which we consider only reactions between gaseous species, with the exception that  $H_2$  and HD forming on the grain surface, and we neglect the three body reactions.

Roberts and Millar (2000a,b) developed new models for the chemistry of deuterium, to investigate the fractionation of doubly-deuterated species, in interstellar molecular clouds. These models depend on a wide range of physical parameters including, density, temperature, elemental abundances and the freeze out of molecules on the dust grains. Our model is partially based on that part of gas-phase reactions of Roberts and Millar (2000a,b) for producing doubly-deuterated species. Our model extended the models of Roberts and Millar (2000a,b) to include multiply-deuterated ammonia.

The mono-deuterated reaction set is complete in the sense that for every reaction containing a hydrogen-bearing molecule there is an analogue reaction containing the equivalent mono-deuterated species. If more than one reaction product contains hydrogen atoms then uncertainly arises as to which will be the deuterium-bearing product. Due to lack of comprehensive experimental data, the usual approach is to assume statistical branching ratio between the various possibilities (Millar et al., 1989; Rodgers and Millar, 1996).

In this study we were interested in the chemistry of doublydeuterated ammonia,  $NHD_2$ . We also interested with multiplydeuterated ammonia, which begin after the formation of  $NH_3$ , with deuterated ions XD<sup>+</sup>. In this model the metals (Fe, Mg, Na and Si) play an important role in determining the ionization fraction.

Since the chemical structure of interstellar clouds depends on the temperature, the radiation field and density number, then we have adopted four models of initial elemental (see Table 1), to study the deuterium chemistry in TMC-1, Orion, L134N and W3 IRS4 clouds. We have neglected the chemistry of species of polycyclic aromatic hydrogen type and their reaction with smaller molecules. As Herbst and Leung (1986) we also neglected the effects of enhanced rate coefficient in ionpolar neutral reactions. The chemical scheme used here is based on that of the most recent UMIST rate file, RATE99 (Le Teuff et al., 2000), with updating the rate coefficients of some modifications by Chastaing et al. (2001). We have adopted the cosmic D/H ratio measured by Linsky et al. (1995). We used a constant ratio for C/O. We neglect the effect of X-ray ionization. The electron abundance is set equal to the sum of the ion abundances.

#### 3. Results and discussion

We have followed a large number of runs for the four models listed in Table 1. Using a different initial elemental abundances and cosmic ray ionization rate, the best initial elemental abundances are given in Table 2.

The cosmic ray ionization rate of  $1.3 \times 10^{-17} \text{ s}^{-1}$  is used for models (1)–(3) and high cosmic ray ionization rate of  $1.3 \times 10^{-16} \text{ s}^{-1}$  is used for model (4). The reduction of the initial elemental abundances and a high cosmic ray ionization rate used in model (4) is in agreement with Amin (1999).

Tables 3–6 compare the molecular D/H ratios observed toward the above four clouds, with the results from our gasphase models at both early  $(10^5 \text{ year})$  time and steady state (100 million year) time, and other theoretical calculations.

Table 1	Our models, where $n$ is number density.						
Model	$n ({ m m}^{-3)}$	$T(\mathbf{K})$	Average (mag.)	Clouds			
1		10	10	TMC-1			
2	1 (4)	70	10	Orion			
3	1 (4)	10	15	L134N			
4	5 (4)	55	30	W3 IRS4			

Table 2Initial fractional abundances.

Species	Model 1	Model 2	Model 3	Model 4
H <sub>2</sub>	0.5	0.5	0.5	0.5
$C^+$	7.3 (-5)	7.3 (-5)	7.3 (-5)	4.0 (-6)
0	1.7 (-4)	1.7 (-4)	1.7 (-4)	8.0(-6)
Ν	2.14 (-5)	2.14 (-5)	2.14 (-5)	5.1 (-6)
S	1.0(-7)	1.0(-7)	1.0(-7)	6.0(-8)
Si	2.0(-8)	2.0(-8)	2.0(-8)	1.5 (-10)
Fe <sup>+</sup>	1.0(-8)	1.0(-8)	1.0(-8)	1.5(-10)
$Mg^+$	1.0(-8)	1.0(-8)	1.0(-8)	1.5 (-10)
Na <sup>+</sup>	1.0(-8)	1.0(-8)	1.0(-8)	1.5(-10)
$H^{3+}$	1.0(-11)	1.0(-11)	1.0(-11)	0.0
HD	1.6(-5)	1.6(-5)	3.2(-5)	5.0 (-6)
He	0.14	0.14	0.14	0.14

Species	Observation	Our cal. M	Our cal. M1			Ref.	
		Early	Steady	Early	Steady		
DCO <sup>+</sup> /HCO <sup>+</sup>	0.02	0.018	0.082	0.019	0.087	Tine et al. (2000)	
NH <sub>2</sub> D/NH <sub>3</sub>	0.009-0.014	0.086	0.029	0.084	0.028	Tine et al. (2000)	
HDCO/H <sub>2</sub> CO	0.0059-0.11	0.043	0.056	0.042	0.055	Turner (2001)	
DCN/HCN	0.023	0.006	0.022	0.009	0.025	Turner (2001)	
DNC/HNC	0.015	0.015	0.015	0.015	0.015	Wootten (1987)	
$C_2D/C_2H$	0.01	0.012	0.028	0.011	0.027	Millar et al. (1989)	
$C_4D/C_4H$	0.004	0.004	0.027	0.004	0.029	Turner (1989)	
$N_2D^+/N_2H^+$	0.08	0.03	0.058	0.025	0.025	Tine et al. (2000)	
$C_3HD/C_3H_2$	0.08-0.16	0.007	0.027	0.006	0.02	Bell et al. (1988)	
$C_3H_3D/C_3H_4$	0.054-0.065	0.082	0.098	0.083	0.099	Gerin et al. (1992)	
DC <sub>3</sub> N/HC <sub>3</sub> N	0.03-0.1	0.008	0.026	0.007	0.026	Howe et al. (1994)	
DC <sub>5</sub> N/HC <sub>5</sub> N	0.013	0.023	0.026	0.023	0.026	MacLeod et al. (1981)	
HDCS/H2CS	0.02	0.04	0.05	0.04	0.046	Minowa et al. (1997)	

 Table 3
 A comparison of abundance ratio measured in TMC-1 cloud with predictions from our model (1) and Roberts and Millar (2000a)

Table 4 A comparison of abundance ratio measured in Orion cloud with predictions from our model (2) and Millar et al. (1989).

Species	Observation	Our cal. M2	M1998	Ref.
DCO <sup>+</sup> /HCO <sup>+</sup>	0.002	0.002	8 (-4)-8 (-5)	Penzias (1979)
NH <sub>2</sub> D/NH <sub>3</sub>	0.003	0.0029	4 (-4)	Walmsley et al. (1987)
HDCO/H <sub>2</sub> CO	0.02	0.019	0.004-0.005	Loren and Wootten (1985)
DCN/HCN	0.006	0.0058	0.001-4 (-4)	Wootten (1987)
DNC/HNC	0.01	0.043	9 (-4)-2 (-4)	Wootten (1987)
$C_2D/C_2H$	0.045	0.032	0.003	Combes et al. (1985), Vrtilek et al. (1985)
CH <sub>3</sub> OD/CH <sub>3</sub> OH	0.01-0.06	0.005	0.003-0.004	Mauersberger et al. (1988)
HDO/H <sub>2</sub> O	> 0.002	0.001	0.001-2 (-4)	Henkel et al. (1987)

Note: a(-b) stands for  $a \times 10^{-b}$ .

Table 5A comparison of abundance ratio measured in L134N (Tine et al., 2000; Roueff et al., 2000) with predictions from our model(3) and Millar (2002).

Species	Observation	Our cal. M2		M2002	
		Early	Steady	Early	Steady
DCO <sup>+</sup> /HCO <sup>+</sup>	0.18	0.026	0.17	0.016	0.047
NH <sub>2</sub> D/NH <sub>3</sub>	0.1	0.006	0.1	0.0075	0.02
NHD <sub>2</sub> /NH <sub>3</sub>	5 (-3)	2.3 (-5)	5.1 (-3)	1.3 (-4)	8.4 (-5)
$N_2D^+/N_2H^+$	0.35	0.03	0.038	0.023	0.039

The most important primary reactions to extract deuterium from HD involve ion-neutral isotope exchange reactions:

Table 6	A co	ompa	arison	of	abu	ndano	ce	ratio	measur	ed in
W3IRS4	cloud	by	Helm	ich	and	van	Di	shock	(1997)	with
prediction	ns fron	1 ou	r mode	el (4	4).					

Species	Observation	Our cal. M2				
HDS/H <sub>2</sub> S	< 9.1 (-2)	1.0 (-3)				
HDCO/H <sub>2</sub> CO	< 3.8 (-2)	0.02				
DCN/HCN	<4.3 (-3)	0.001				
DNC/HNC	7.1 (-3)	9 (-4)				
$C_2D/C_2H$	0.045	0.032				
CH <sub>3</sub> OD/CH <sub>3</sub> OH	< 7.1 (-2)	0.006				
Note: $a(-b)$ stands for $a \times 10^{-b}$ .						

$\mathrm{H}_{2}^{+} + \mathrm{HD} \to \mathrm{HD}^{+} + \mathrm{H}_{2}, \tag{1}$	1)	ļ
	- /	1

 $CH_3^+ + HD \rightarrow CH_2D^+ + H_2, \qquad (2)$ 

 $C_2H_2^+ + HD \rightarrow C_2HD^+ + H_2, \tag{3}$ 

where the reaction exoergicities although small  $-\Delta E_1/k = 220 \text{ K}, \Delta E_2/k = 375 \text{ K}, \Delta E_3/k = 550 \text{ K}$  – are much larger than the temperatures of cold interstellar clouds. At low temperatures, the reverse reactions do not occur efficiently despite the large abundance of H<sub>2</sub>. Once formed these deuterated ions can pass on their enhanced deuterium content to other species in chemical reactions (Millar, 2002). In addition to the forward and reverse reaction in (1), H<sub>2</sub>D<sup>+</sup> can be destroyed by metals, by dissociative recombination with electron,

$$\mathbf{H}_{2\mathbf{D}}^{+} + \mathbf{e} \to \mathbf{H} + \mathbf{H} + \mathbf{D}, \tag{4a}$$

 $H_2D^+ + e \to H_2 + D, \tag{4b}$ 

 $H_2 D^+ + e \rightarrow H + H D \tag{4c}$ 

(14)

with total rate coefficient equal to  $6 \times 10^{-8} (T/300)^{-0.5} \text{ cm}^{-3} \text{ s}^{-1}$  (Larsson et al., 1996) and by reaction with neutral molecules (CO, N<sub>2</sub>, H<sub>2</sub>D).

The reactions for  $CH_2D^+$  and  $C_2HD^+$  are similar to that of  $H_2D^+$ , except that each of these ions undergo a radiative association reaction with  $H_2$  (Millar et al., 1989; Herbst et al., 1987)

$$CH_2D^+ + H_2 \rightarrow CH_4D^+ + photon,$$
 (5)

$$C_2HD^+ + H_2 \rightarrow C_2H_3D^+ + photon.$$
(6)

Our calculated ratios for  $H_2D^+/H_3$ ,  $CH_2D^+/CH_3^+$  and  $C_2HD^+/C_2H_2^+$  are all enhanced at low temperature, but  $H_2D^+$  is responsible for D/H ratio other than the two molecules. At high temperature  $H_2D^+$  is rapidly destroyed by  $H_2$ , so  $CH_2D^+$  and  $C_2HD^+$  are responsible for D/H ratio.

Our calculated ratio of  $C_2D/C_2H$  is in agreement with observations of TMC-1 at an early (10<sup>5</sup> year) time and higher at the steady state (10<sup>8</sup> year) time. This result is the same as that obtained by Roberts and Millar (2000a). The C<sub>2</sub>D species is formed by:

$$C_2HD^+ + e \to C_2D + H, \tag{7}$$

$$C_2H + D \rightarrow C_2D + H + 580 \text{ K}$$
(8)

and is destroyed by

$$C_2 D + H \to C_2 H + D. \tag{9}$$

In model (2) with temperature of 70 K, the neutral–neutral rate coefficient of reaction (8) and dissociative recombination of  $C_2HD^+$ , become more competitive, and fractional abundance of  $C_2D$  increases. This result agrees with that of Herbst et al. (1987). Also  $C_2D$  is formed from cosmic ray induced photodissociation of  $C_2HD$ . At steady state time our calculated  $C_2D/C_2H$  ratio is in agreement with observations of Orion cloud and greater than that obtained by Millar et al. (1989).

As Millar et al. (1989)  $C_2H_3D^+$  is precursor to  $C_2H_2D$ , so the reactions of  $C^+$  with  $C_2H_2D$  transfer fractionation to  $C_3$ -bearing hydrocarbons via

$$C_2H_2D + C^+ \to C_3HD^+ + H \tag{10}$$

and subsequently to  $C_3D$  via dissociative recombination. The  $C_3H_2$  molecule is not transferred easily to the  $C_3HD^+$  and  $C_3HD$  is formed as (see Bell et al., 1988);

 $C_2HD + C^+ \rightarrow C_3D^+ + H, \qquad (11)$ 

 $C_3D^+ + H_2 \rightarrow C_3H_2D^+ + photon, \qquad (12)$ 

$$C_3H_2D^+ + e \to C_3HD + H.$$
(13)

Our calculated  $C_3HD/C_3H_2$  ratios in both an early and the steady state times are not in agreement with observation of TMC-1 cloud.

In our model C +  $3H_3D$  molecule is formed through the  $CH_2D^+$  ion, by dissociative recombination of  $CH_2DC_2H_2^+$  and  $CH_2DC_3H_2^+$  ions, which are themselves formed from  $CH_2D^+$  by ion-neutral reactions with small hydrocarbon species like methane or acetylene.  $C_3H_3D$  is destroyed by atomic and molecular ions, primarily  $H_3^+$  and  $He^+$ . Our calculated  $C_3H_3D/C_3H_4$  ratios at an early and the steady state time are greater than the lower and upper limit of observations of TMC-1 cloud. This results is the same as that of Roberts and Millar (2000a).

In our model  $C_4D$  species is formed by the following reactions:

$$C_4 + HD \rightarrow C_4D^+ + H, \qquad (14)$$

$$C_4D^+ + H_2 \rightarrow C_4HD^+ + H, \tag{15}$$

$$C_{3}HD^{+} + C^{-} + C_{4}HD^{-} + H, \qquad (10)$$

$$C_4\Pi D^+ + e \to C_4D + \Pi \tag{17}$$

and is destroyed by

 $C^+$  + UD + C D<sup>+</sup> + U

$$C_4 D + O \rightarrow C_3 D + CO, \tag{18a}$$

$$C_4D + H_3^+ \rightarrow C_4HD^+ + H_2, \qquad (18b)$$

$$C_4 D + HCO^+ \rightarrow C_4 HD^+ + CO, \qquad (18c)$$

$$7C_4D + H_3O^+ \rightarrow C_4HD + H_2. \tag{18d}$$

At an early time our calculated  $C_4D/C_4H$  ratio is in agreement with observations of TMC-1 clouds.

The D–N bond begins due to the lower proton (deuteron) affinity of  $H_2D^+$ , which reacts with  $N_2$  to form  $N_2D^+$  species,

$$H_2D^+ + N_2 \rightarrow N_2D^+ + H_2.$$
 (19)

In addition to reaction (19), D atom can also react with  $N_2H^+$  to form  $N_2D^+$ 

$$N_2H^+ + D \rightarrow N_2D^+ + H.$$
<sup>(20)</sup>

Reaction (20) has an exothermicity of about 550 K (Adams and Smith, 1985). Due to the proton affinity of  $N_2$  is very small, so that  $N_2D^+$  reacts with  $CH_2D^+$  and  $C_2HD^+$ .

IN our model  $N_2D^+/N_2H^+$  ratios are less than the observations of TMC-1 and Orion clouds. This is because at dense clouds  $N_2D^+$  condense onto grains, i.e.  $N_2 D^+$  is not easily predictable by gas-phase reactions and its value must be predicted through evaporation from the grain (Willacy and Millar, 1998; Millar et al., 1989; Millar, 2002).

After  $NH_3$  is formed by the reaction sequence  $N_2 \xrightarrow{He_+} N^+ \xrightarrow{4H_2} NH_4^+ \xrightarrow{e} NH_3$ , deuteron transfer reaction forms  $NH_3D^+$  which can then recombine to give  $NH_2D$  as;

$$NH_3 + XD^+ + \rightarrow NH_3D^+ + XH, \tag{21}$$

$$NH_3^+ + e \to NH_2D + H, \tag{22}$$

where  $XD^+$  represents all species capable of transferring a proton or deuteron to NH<sub>3</sub>, principally H<sub>3</sub><sup>+</sup>, N<sub>2</sub>H<sup>+</sup>, HCO<sup>+</sup> and their deuterated isotopomers. Successive deuteron transfer reaction can lead eventually to NHD<sub>2</sub> and ND<sub>3</sub> as

$$NH_2D + XD^+ \rightarrow NH_2D^+ + +X, \tag{23}$$

$$NH_2D_2^+ + e \to NHD_2 + H, \tag{24}$$

$$\mathrm{NHD}_2 + \mathrm{XD}^+ \to \mathrm{NHD}_3^+ + \mathrm{X}, \tag{25}$$

$$\mathrm{NHD}_3^+ + \mathrm{e} \to \mathrm{ND}_3 + \mathrm{H},\tag{26}$$

$$\mathrm{NHD}_3^+ + \mathrm{e} \to \mathrm{NHD}_2 + \mathrm{D}. \tag{27}$$

From the last reactions (21)–(27) the relative fractional abundances depend on the  $XD^+/XH^+$  and the branching ratio for dissociative recombination of the deuterated ions.

By assuming the rate coefficients of reactions (26) and (27) are equal, our calculated fractional abundances for ND<sub>3</sub> is about  $2 \times 10^{-11}$ . This value is greater than that obtained by Rodgers and Charnely (2001) by two times. The abundance of ND<sub>3</sub> in our model is in the same order of magnitude as given by Lis et al. (2002) for B1 cloud. From this result we can conclude that ND<sub>3</sub> can be detected in L134N cloud.

One species for which time dependence significant is  $DCO^+$ . The species  $DCO^+$  is more readily observable than  $N_2D^+$ , because the abundance of CO is greater than that of

 $N_2$  in interstellar clouds. So at low temperature the D–C band begins with the reaction:

$$H_2D^+ + CO \rightarrow DCO^+ + H_2 \tag{28}$$

at an early time. Also DCO<sup>+</sup> is formed by:

$$C^{+} + HDO \rightarrow DCO^{+} + H, \tag{29}$$

$$HCO^+ + D \rightarrow DCO^+ + H.$$
 (30)

At the low temperature and density, the atomic deuterium is very abundant, so reaction (30) proceeds very rapidly and can further enhance the fractionation of  $DCO^+$ .

At high temperature the abundance of  $CH_4D^+$  is greater than  $H_2D^+$ , then  $DCO^+$  is formed by,

$$CH_4D^+ + CO \rightarrow DCO^+ + CH_4.$$
 (31)

In addition to this reaction, there are significant contributions to  $DCO^+$  formation in the reactions

 $CH_2D^+ + O \rightarrow DCO^+ + H_2, \tag{32}$ 

$$C_2 HD^+ + O \to DCO^+ + CH, \tag{33}$$

$$CD + O \rightarrow DCO^+ + e$$
,

$$CH_2D^+ + O \rightarrow DCO^+ + H_2. \tag{35}$$

The primary of D in dark clouds is the dissociative recombination of DCO<sup>+</sup>,

$$DCO^+ + e \rightarrow CO + D.$$
 (36)

Our calculated  $DCO^+/HCO^+$  ratio is in agreement with observations of TMC-1 cloud at an early time. At the steady state time it is in agreement with observations of Orion and L134N clouds. This result differs than that obtained by Millar et al. (1989) and Millar (2002).

The band between D–C–O begins by the formation of deuterated formaldehyde, HDCO, which is formed from two species derived from  $CH_2D^+$ ,  $CH_4D^+$  and  $CH_2D$ , via the reactions;

$$CH_4D^+ + O \rightarrow H_2DCO^+ + H_2 \tag{37}$$

followed by dissociative recombination

$$H_2DCO^+ + e \rightarrow HDCO + H$$
 (38)

and

$$CH_2D^+ + O \rightarrow HDCO + H^+.$$
 (39)

HDCO also reacts with  $H_3^+$  to form  $H_2DCO^+$  which will then recombine to HDCO. The calculated HDCO/H<sub>2</sub>CO ratio is in agreement with observations of TMC-1, Orion and W3IRS4 clouds.

Also the band between D–C–N begins by the formation of DCN, which is formed a derivative of  $CH_2D^+$ ;

$$CH_2D^+ + e \rightarrow CHD + H,$$
 (40)

$$CHD^+ + N \to DCN + H^+.$$
<sup>(41)</sup>

The main derives of DCN at high temperature is the neutralneutral reactions;

$$DCO^{+} + N \rightarrow DCN + O, \tag{42}$$

$$HCN^{+} + D \rightarrow DCN + H.$$
<sup>(43)</sup>

We found that the reaction

$$H^{+} + DCN \rightarrow HCN^{+} + D \tag{44}$$

can cycle deuterium between the atomic D and the DCN molecule. This result is in agreement with that of Schilke et al. (1992) and Hatchell et al. (1998). As shown in Tables 3, 4 and 6 our calculated DCN/HCN is in agreement with observations of TMC-1, Orion and W3IRS4 clouds.

At the low temperature, the dominant route to deuterated cyanoacetylene,  $DC_3N$ , formation is thought to be

$$C_3H_2D^+ + N \to HDC_3N^+ + H \tag{45}$$

followed by

$$HDC_3N^+ + e \to DC_3N + H. \tag{46}$$

At the high temperature, the deuterated acetylene should form deuterated cyanoacetylene through

$$C_2HD + CN \to DC_3N + H. \tag{47}$$

 $DC_3N$  species is only observed in TMC-1 cloud, and our calculated its ratio is in agreement with observations at steady state time.

In our models we have assumed that the species  $CH_3OD$  is formed from the radiative association reaction

$$CH_3^+ + HDO \rightarrow CH_3OHD^+ + photon$$
 (48)

and

(34)

$$CH_3OHD^+ + e \rightarrow CH_3OD + H.$$
 (49)

The HDO species comes from the rapidly exothermic reaction

$$CH_2D^+ + H_2O \to CH_3 + HDO.$$
<sup>(50)</sup>

The calculated CH + 3OD/CH + 3OH ratio in our model differs from observations. This is because, at high densities most molecules condense onto grain; i.e. CH<sub>3</sub>OH and CH<sub>3</sub>OD is not easily predictable by gas-phase reaction and their values must be predicted through evaporation from the grains surface (Amin, 1999). Also our calculated ratio for HDO/H<sub>2</sub>O is small, because the large value must be predicted by shock chemistry (Amin, 2001; Parise et al., 2002).

When we compare our predicted gas-phase abundances with those observed in TMC-1 and W3IRAS4 clouds, only HDCS is in agreement and other species HDS is smaller than observation. In our model after the formation of  $H_2S$  by the chain;

$$\mathbf{S^+} \stackrel{\mathrm{H2}}{\to} \mathbf{HS^+} \stackrel{\mathrm{H2}}{\to} \mathbf{H_3S^+} \stackrel{\mathrm{e}}{\to} \mathbf{H_2S},$$

(By the last chain the calculated fractional abundance of  $H_2S$  is small) deuteron transfer reaction forms  $H_2DS^+$  which then recombine to give HDS as:

$$H_2D^+ + H_2S \rightarrow H_2DS^+ + H_2, \tag{51}$$

$$H_2DS^+ + e \rightarrow HDS + H.$$
 (52)

Also DCO<sup>+</sup> and 
$$H_2D^+$$
 react with  $H_2CS$  to form HDCS as;

$$DCO^+ + H_2CS \rightarrow H_2DCS^+ + CO,$$
 (53)

$$H_2D^+ + H_2CS \rightarrow H_2DCS^+ + CO, \tag{54}$$

$$H_2DCS^+ + e \rightarrow HDCS + H.$$
 (55)

#### 4. Conclusions

With a more extensive chemical network, we made a detailed study of a pseudo-time dependent chemical evolution of deuterium species in different interstellar clouds, TMC-1, Orion, L134N and W3IRS4, with different densities and temperatures. This has been done using different initial elemental abundances and without the temperature dependence of the ion-dipole molecule collisions. We have shown that large abundances of NH<sub>2</sub>D and NHD<sub>2</sub> can be produced by gas phase chemistry in cold dense clouds. Ammonia is deuterated via deuteron transfer from species such as  $H_2D^+$ ,  $DCO^+$  and  $N_2D^+$ , followed by dissociative recombination. We predict the abundance of ND<sub>3</sub> is  $2 \times 10^{-11}$ , by a assuming the rate coefficients of reactions NHD<sub>3</sub><sup>+</sup> + e  $\rightarrow$  ND<sub>3</sub> + H, and NHD<sub>3</sub><sup>+</sup> + e  $\rightarrow$  $NHD_2 + D$  are equal. So we suggest that triply-deuterated ammonia could be detectable in L134N cloud. We have included the fractionation of sulphur-bearing molecules and found a good agreement with observation for HDCS. The very slow formation rates of HDS in cold gas make this molecule particularly useful in probing regions where grain surface chemistry may be important. Reduction of the initial elemental abundances and high cosmic ray ionization gave us good relative abundances for most of the observed deuterated species in W3IRS4.

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#### Further reading

Brown, R.D., Rice, E.H.N., 1981. Interstellar deuterium chemistry. Philosophical Transactions 303A (1480), 523–533. احتراق جزيئات الديوتريوم في السحب البين نجمية الكثيفة الباردة والدافئة

مجدي يوسف أمين و أيمن سعيد كردي قسم الفيزياء والفاك ، كلية العلوم، جامعة الملك سعود ص ب 2455 ، الرياض 11451

(قدم للنشر في 1429/2/24 هـ؛ وقبل للنشر في 1/1/24 هـ)

ملخص البحث. حيث أن جزيئات الديوتريوم تستخدم في: 1- اختبار القوانين الفيزيائية لسحب ما بين النجوم. 2- در اسة علاقة الربط بين تكوين الثلج في مادة ما بين النجوم والمذنبات. 3- طريقة تكوين نظائر المركبات في مادة ما بين النجوم لهذا تقوم در اسات نظرية وعملية (عن طريق رصد السحب البين نجمية بطرق مختلفة)، لفهم عملية تكوين

وتكوين مركبات الديوتريوم وحساب كمية الوفرة في سم<sup>2</sup> .

في هذا البحث نقوم بدر اسة طريقة تكوين مركبات الديوتريوم وحساب كمية الوفرة في أربعة سحب بين نجمية ذات عوامل فيزيائية مختلفة وهي TMC-1 و Orion و L134N و W3IRS4 وفي هذه الدراسة تم تسيد برنامج كيميائي يحتوي على 5340 تفاعل كيميائي لـ 0014 مركب وعنصر كيميائي منهم 130 مركب للديوتريوم. وقد تم تحويل كيميائي يحتوي على 5340 تفاعل كيميائي لـ 608 مركب وعنصر كيميائي منهم 130 مركب للديوتريوم. وقد تم تحويل هذه التفاعلات إلى معادلات تفاعل كيميائي لـ 608 مركب وعنصر كيميائي منهم 130 مركب للديوتريوم. وقد تم تحويل هذه التفاعلات إلى معادلات تفاصلية من الدرجة الأولى تحقق شرط Stiff وحلها عدديا باستخدام طريقة الـ Gear. وقد تم الحصول على قيم متوافقة مع الأرصاد الفلكية باستخدام قيم بدائية مختلفة للعناصر الأساسية حسب الوضع الفيزيائي الحصول على قيم متوافقة مع الأرصاد الفلكية باستخدام قيم بدائية مختلفة للعناصر الأساسية حسب الوضع الفيزيائي للسحابة المدروسة. وقد تم الفيزيائي منام NH20 و عنصر كيميائي منهم 300 مركب قدم تحويل هذه التفاعلات إلى معادلات تفاضلية من الدرجة الأولى تحقق شرط NH2f وحلها عدديا باستخدام طريقة الـ Gear. وقد تم الحصول على قيم متوافقة مع الأرصاد الفلكية باستخدام قيم بدائية مختلفة للعناصر الأساسية حسب الوضع الفيزيائي الحصول على قيم وفرة كبيرة لكل من NH2f و NHD2 في السحب الكثيفة الباردة عند حصول حللة الثبات في عملية التكوين والتكسير عند زمن قدرة مائة مليون سنة.