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# Investigation on the key features of L-Histidinium 2-nitrobenzoate (LH2NB) for optoelectronic applications: A comparative study



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#### **KEYWORDS**

Organic compounds; Raman spectroscopy and scattering; Optical materials; Nonlinear optical material; Computational techniques **Abstract** The current work is to highlight the fundamental acumen about the molecular structure, photophysical and static first hyperpolarizability ( $\beta$ ) of L-Histidinium 2-nitrobenzoate (LH2NB) organic molecule for the first time. Hartree–Fock (HF) and density functional theory (DFT) has been applied using different functional at 6-31G<sup>\*\*</sup> basis set for the first time. The strong correlation has been observed between experimental and theoretical vibrational spectra. TD-DFT method has been used at different levels of theory to study the UV–Visible spectra. The analysis of HOMO and LUMO was done to explain the charge interaction taking place within the molecule and the energy gap was evaluated. The value of dipole moment is found to be lower in excited state than ground state as calculated from all applied methods. The value of total static first hyperpolarizability was found to be 7.447 × 10<sup>-30</sup> esu at B3LYP/6-31G<sup>\*\*</sup> level of theory, which is about 20 times higher than urea molecule. The current results indicate that the studied molecule may be a decent applicant for opto-electronic applications.

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#### 1. Introduction

Since last few decades the research and development on L-Histidine and its complexes has been receiving an immense

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attention due to their easy synthesis, growth and good nonlinear optical (NLO) properties and hence emerges as one of the most extensively explored amino acid from its family. It can forms a variety of complexes with different organic and inorganic materials such as: gold (III)-L-histidine (Cuadrado et al., 2000), L-Histidinium 2-nitrobenzoate (Moovendaran et al., 2013; Natarajan et al., 2012), l-histidine-4-nitrophenolate 4-nitrophenol (LHPP) Dhanalakshmi et al., 2010, 1-histidine acetate (Madhavan et al., 2007), L-Histidinium perchlorate (LHPCL) Aruna et al., 2007, 1-Histidine nitrate (Dhas and Natarajan, 2008), L-histidinium trifluoroacetate (Dhas et al., 2008), 1-histidine hydrofluoride dihydrate (LHHF) Madhavan et al., 2006, 1-histidine hydrochloride monohydrate (Anandan and Jayavel, 2011; Madhavan et al., 2007), and 1-histidine

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bromide (Ahmed et al., 2008) and shows noticeable Second Harmonic Generation (SHG) efficiency. L-Histidine has another important advantage of being an organic material, which is well known due to its low cost, extraordinary nonlinearity, high optical threshold, synthetic litheness, and easy molecular design. Also its configuration can be modified to get the desired nonlinear optical (NLO) properties for tailor made applications and also shows low dielectric constants which makes it useful in terahertz (THz) generation devices (Boomadevi et al., 2004; Moovendaran et al., 2012; Zyss et al., 1984; Chemla, 2012; Shakir et al., 2010, 2014; Ledoux et al., 1987; Fujiwara et al., 2006; Shakir et al., 2009).

The synthesis of a new L-Histidine compound named L-Histidinium 2-nitrobenzoate (LH2NB) (chemical structure shows in Fig. 1) has been reported (Natarajan et al., 2012), and its crystallization, molecular structure, vibrational, optical, second harmonic generation (SHG) efficiency and thermal properties are described. Moovendaran et al. (2013) has reported the SHG efficiency of titled compound which is about 2 times higher than standard KDP crystal.

As per the current available literature it is clear that only experimental studies have been performed on the titled compound so far. However, it is very important and justified to study its theoretical properties such as geometrical, vibrational, photophysical, nonlinear etc. using HF and density functional theory (DFT) to understand the mechanism responsible for its use in various optoelectronic applications. Hartree-Fock (HF) and DFT show key advantages in calculating the various parameters such as: vibrational frequencies, molecular geometries of different kinds of materials very precisely within short time and at low cost (Johnson et al., 1993; Cinar et al., 2014; Shkir and Abbas, 2014a, 2014b; Arivazhagan and Meenakshi, 2012; Reshak and Khan, 2014; Govindarasu and Kavitha, 2014; Elleuch et al., 2007; Shkir et al., 2014, 2015a, 2015b, 2015c, 2015d). Furthermore, the range separated functionals such as CAM-B3LYP, wb97xd and many more are efficient to calculate the electronic and nonlinear optical (NLO) properties which are much superior to the conventional methods (Johnson et al., 1993; Cinar et al., 2014; Shkir and Abbas, 2014a, 2014b; Arivazhagan and Meenakshi, 2012; Reshak and Khan, 2014; Govindarasu and Kavitha, 2014; Elleuch et al., 2007; Shkir et al., 2014, 2015a, 2015b, 2015c; Abbas et al., 2015).

In the current work, the author's goal is to highlight the key features of LH2NB molecule by HF and DFT (using different functional) studies carried out for the first time. It may be noted here that the reason for applying the different functional is to have a in-depth knowledge about the appropriateness of functional which gives better results for the titled molecule as every functional has different extension of DFT.



Figure 1 Chemical structure of LH2NB.

#### 2. Computational details

HF (Fischer et al., 1973) and DFT using B3LYP- Becke's three parameter exchange functional B3 combined with Lee-Yang-Parr correlation functional LYP (Becke, 1993; Lee et al., 1988) for obtaining the molecular structures, IR and Raman spectra. Further the TD-DFT study has been performed using B3LYP along with range separated functional such as CAM-B3LYP (Yanai et al., 2004), wb97xd (Chai and Head-Gordon, 2008), PBE0 (Adamo and Barone, 1999), M06 (Zhao and Truhlar, 2008) for calculating opto-electronic properties (Dreuw and Head-Gordon, 2004; Foster and Wong, 2012; Wong et al., 2009; Gibbs et al., 2011). The stable geometry was achieved following the true minimum on the potential energy surface (PES) attained by solving the self-consistent field equation. Infrared (IR) and vibrational (Raman) frequencies were calculated using optimized structural parameters to characterize all stationary points as minima. All the theoretical calculations were made using Gaussian 09W program package (Frisch et al., 2009) with the default convergence principles, without any constraint on the geometry. By applying the different functional the dipole moment, polarizability, anisotropy of polarizability, and static and total first hyperpolarizability values were calculated. Finite Field (FF) method was employed to calculate the value of total first hyperpolarizability ( $\beta_{tot}$ ) and its tensor components. FF method was generally applied to know the nonlinear optical properties because this approach can be used in concert with the electronic structure method to work out  $\beta$  values.  $\beta_{tot}$  values calculated by this method is found to be genuine with experimental structure property relationship recently. A static electric field (F) has been applied to a molecule in FF method and the energy (E) is expressed by the following relation:

$$E = E^{(0)} - \mu_1 F_1 - \frac{1}{2} \alpha_{ij} F_i F_j - \frac{1}{6} \beta_{ijk} F_i F_j F_k - \frac{1}{24} \gamma_{ijkl} F_i F_j F_k F_l - \dots$$
(1)

where  $E^{(0)}$  is the energy of molecule in the absence of an electric field,  $\mu$  is components of the dipole moment vector,  $\alpha$  is the linear polarizability tensor,  $\beta$  and  $\gamma$  are the first and second hyperpolarizability tensors respectively, while i,j and k label the x,y and z components respectively. Values of  $\mu$ ,  $\alpha$ ,  $\beta$ , and  $\gamma$  can be obtained by differentiating Eq. (1) with respect to *F*.

The value of static hyperpolarizability ( $\beta_0$ ) can be calculated from the following equation:

$$\beta_0 = \frac{3}{5} \beta_{tot} \tag{2}$$

Further, the optical absorption spectra were calculated by time dependent DFT (TDDFT) study suing different functional.

GCRD parameters of the titled molecule have been calculated as follows:

A relation for absolute hardness ( $\eta$ ) was established Parr and Chattaraj, 1991; Pearson, 1985; Parr and Pearson, 1983 i.e.:

$$\eta = \frac{I - A}{2} \tag{3}$$

where I is the vertical ionization potential energy and A is vertical electron affinity.

Koopman's theorem is associated within the structure of HF self-consistent field molecular orbital theory (Koopmans, 1933), the ionization energy and electron affinity can be specified over HOMO and LUMO orbital energies as given below:

$$I = -E_{\text{HOMO}}$$
 and  $A = -E_{\text{LUMO}}$ 

*I* and *A* values of LH2NB molecule are presented in Table 5. Greater HOMO energy is related to the more reactive molecule in the reactions with electrophile, while minor LUMO energy is necessary for molecular reactions with nucleophile (Rauk, 2001). The hardness of any molecule is related to the HOMO–LUMO energy gap and expressed as.

$$\eta = \frac{1}{2} (E_{\text{LUMO}} - E_{\text{HOMO}}) \tag{4}$$

Electronic chemical potential is achieved by:

$$\mu = -\left(\frac{I+A}{2}\right) \tag{5}$$

Chemical softness is achieved by:

$$\mathbf{S} = \frac{I}{2\eta} \tag{6}$$

Electronegativity is achieved by:

$$\chi = \left(\frac{I+A}{2}\right) \tag{7}$$

The electrophilicity index is achieved by:

$$\omega = \frac{\mu^2}{2\eta} \tag{8}$$

#### 3. Results and discussion

#### 3.1. Molecular geometry

The stable molecular geometry of LH2NB was achieved by HF and B3LYP using 6-31G\*\* basis set as shown in Fig. 2 (a) and (b) respectively. The coordinates used in the current work for theoretical calculations were earlier reported CCDC-857702 (Natarajan et al., 2012). The geometry was also optimized by other methods such as B3LYP, CAM-B3LYP and wb97xd levels of theory using 6-31G\*\* basis set. Some important geometrical parameters such as bond lengths, bond angles are tabulated in Table 1 calculated by HF and B3LYP and compared with experimental values (Natarajan et al., 2012) and found in good relationship. The hydrogen bonding is denoted by peppered line and the values of these hydrogen bonds are given in Table 2 obtained at all levels of theory. As clear from Table 2 the geometry optimized by HF is found to be in better agreement than other methods with experimental (Natarajan et al., 2012).

Therefore, here author focused on geometrical parameters obtained from the HF method, as clear from figure (Fig. 2) that several hydrogen bondings have been observed between L-Histidinium (act as cation) and p-nitrobenzoic acid (act as anion). Three intermolecular hydrogen bondings such as the first one between H(17) and O(35) atoms i.e. N(12)–H(17)—O(35) [H(17)—O(35) = 1.7215 Å], second between H(20) and O(34) atoms i.e. N(14)–H(20)—O(34) [H(20)—O(34) = 1.6902 Å], and third one between H(11) and O(36) atoms

i.e. C(10)-H(11)-O(36) [H(11)-O(36) = 2.6428 Å] have been observed. Three intramolecular hydrogen bondings were observed in L-histidine molecule itself between N(12) and (H (18)atoms i.e. N(12)-H(18)-O(15) [H(18)-O(15) = 2.9384 Å], N(12) and H(19) atoms i.e. N(12)-H(19)-O (15) [H(19)-O(15) = 1.9553 Å] and H(9) and O(16) atoms i.e. C(8)-H(9)-O(16) [H(9)-O(16) = 2.6788 Å], respectively. The bondings observed by HF as well as by other methods are given in Table 2 along with experimentally reported values (Natarajan et al., 2012) and it shows that molecular geometry obtained by HF is in good agreement than other methods. It may be mentioned here that by other applied methods only four bondings were observed. Because of a large number of hydrogen bondings and charge transfer in the titled molecule it is expected that it will show great NLO properties (Cole et al., 2001). The reason for the difference in bond lengths and angles is that the experimental values were obtained from the X-ray diffraction of the crystalline material in solid crystal form, though the geometry optimization of LH2NB was performed for an isolated molecule. The lowest value of  $C_7$ - $C_8$ (1.343 Å) less than the standard value (1.54 Å) of C–C bond length was observed. Correspondingly, the calculated shortest value for C–O bond length comes out to be 1.222 Å, which is less than the standard value (1.43 Å) of C-O bond length. The other bond lengths such as C-C, C-O, C-N, N-O, and C-H etc. in LH2NB are inside the range of typical values. The bond lengths of C-H are remain between 1.091 Å and 1.10 Å (Shkir and Abbas, 2014b).

#### 3.2. Vibrational analysis

It is well known in the literature that the vibrational (Infrared and Raman) spectroscopic techniques have been widely applied by the organic chemists to study the functional groups, bonding to different molecular conformations and reaction mechanisms by tentatively assigning their observed fundamental modes (Teimouri et al., 2009; Colthup, 2012; Socrates, 2004; Dollish et al., 1974; Smith, 1998; Roeges, 1994). It is well known that when hydrogen (H), nitrogen (N), oxygen (O) etc. atoms are present between two molecules or within a molecule the inter-molecular and intra-molecular hydrogen bonding appears. The calculated IR and Raman spectra of LH2NB using HF and B3LYP methods are shown in Fig. 3. The theoretically calculated frequencies can be well matched with the experimentally observed frequencies by applying the scaling factor (for HF, 0.8929 and B3LYP, 0.9613) (Sinha et al., 2004; Alcolea Palafox, 2000; Shkir et al., 2015). IR and Raman frequencies of various theoretically predicted and experimentally observed peaks with their tentative assignment are listed in Table 3 and comparison exhibits strong agreement with the reported values.

The vibrational modes present in the molecule at  $\sim$ 3177, 3187 cm<sup>-1</sup> in HF and B3LYP respectively, are due to NH symmetric stretching vibration. In the similar manner the band observed at 3105, 3113 cm<sup>-1</sup>, and 3153 cm<sup>-1</sup> have been tentatively assigned to NH<sub>3</sub> stretching vibration mode. The CH<sub>2</sub>, NH<sub>3</sub> symmetric stretching vibrations are observed in the region of 2800–3001 cm<sup>-1</sup>. The band at 2478, 2470 cm<sup>-1</sup>-observed in Raman not in IR may be assigned to C–H combinational overtone. The COO<sup>-</sup> ions of carboxylic group shows asymmetric and symmetric stretching characteristic modes at



Figure 2 Optimized stable molecular geometry of LH2NB by (a) HF and (b) B3LYP using 6-31G<sup>\*\*</sup> basis set.

1650, 1634, 1639 cm<sup>-1</sup> and 1433, 1440 cm<sup>-1</sup> in IR and Raman respectively. COO<sup>-</sup> bending and wagging vibration modes were observed at 815, 799 cm<sup>-1</sup> and 823, 808 cm<sup>-1</sup> in IR transmittance and Raman spectra respectively. The band observed at 1393, 1361 cm<sup>-1</sup> and 1353 cm<sup>-1</sup> have been assigned to CH<sub>2</sub> deformation. The asymmetric and symmetric stretching modes of NO<sub>2</sub> functional group is expected in the region of 1370–1330 cm<sup>-1</sup> and observed at 1361, 1313 cm<sup>-1</sup> and 1330, 1336 cm<sup>-1</sup> in these IR transmittance and Raman spectra respectively. The wagging and rocking modes of vibration of NO<sub>2</sub> group are observed at 694 cm<sup>-1</sup> and 690 cm<sup>-1</sup>. The other vibration modes observed in LH2NB molecule are tentatively assigned in Table 3 which exhibits very well correlation between theoretically calculated and experimentally observed peaks in their respective IR and Raman spectra.

#### 3.3. Optical (TD-DFT) study

Time dependent DFT (TD-DFT) is one of the broad approaches in quantum chemistry as well as in solid state physics to calculate the excited state electronic structure. The contemporary DFT methods display a promising stability between accuracy and computational efficiency, when equated to the traditional ab initio and semi-empirical approaches. To examine the kind of electronic transition in LH2NB molecule in gas phase, the TD-DFT has been applied at different functional. As almost precise absorption wavelength can be easily detected at relatively small computing time by such study on the basis of optimized ground state geometry, which is related to the vertical electronic transitions (Kostova et al., 2010: Jacquemin et al., 2005, 2004; Cossi and Barone, 2001). The UV-Vis spectrum was also calculated for the optimized geometry at TD-HF level of theory using 6-31G<sup>\*\*</sup> basis set as shown in Figure 1S (see supplementary data). Absorption wavelength, excitation energies, oscillator strengths and dipole moments were evaluated at the ground state of the optimized geometries using TD-B3LYP, TD-CAM-B3LYP, TD-wb97xd, TD-PBE0, TD-M06 levels of theory using 6-31G\*\* basis set. These calculated values are presented in Tables 4 and 5. The calculated UV-Vis spectrum of LH2NB has only one absorption band at all applied functional as shown in Fig. 4. The value of excitation wavelength is found to be  $\sim$ 328.19 (3.778 eV), 309.07 (4.012 eV), 310.64 (3.990 eV) nm, 317 (3.935 eV) nm and 324

(3.820 eV) nm calculated at B3LYP, CAM-B3LYP, wb97xd, PBE0 and MO6 levels of theory using 6-31G<sup>\*\*</sup> basis set, respectively. The oscillatory strength describes the strength of optical or molecular interactions. The value of f for the titled molecule is found to be between 0.02 to 0.05 obtained at all applied methods. The value of f is used to compare a quantum mechanical transition to that expected by a fully allowed set of classical electromagnetic oscillators. The value of  $f = \sim 1$ , represents a strong transition while generally a quantum mechanically forbidden transition may have f = 0.001. It may be mentioned here that the optimized coordinate at B3LYP has been used in PBE0 and MO6. In experimentally reported result on optical transmission there is a absorption band at about 300 nm (Moovendaran et al., 2013), which is in good agreement with the theoretical value calculated at almost all levels of theory, however it is more close to CAM-B3LYP. The present study reveals that the titled compound owns the property of extended ultraviolet-visible transmittance (Govindarajan et al., 2011; Srinivasan et al., 2006). The obtained high excitations energy (or optical absorption spectra which help us to determine the optical transparency and band gap of the materials) is maybe useful in various optoelectronic devices (Shkir et al., 2014; Kushwaha et al., 2011; Shakir et al., 2009; Kushwaha et al., 2014). The optical band gap was calculated by using Tauc's relation from the absorption spectra obtained from B3LYP and found to be about 3.3 eV. This shows that the titled material possesses high band gap in comparison to other optical materials as well as comparable (Shkir et al., 2014; Babu et al., 2010; Klimm, 2014) and can be used in optical devices. The change in the dipole moment (Table 5) upon excitation was observed which may be due to charge redistribution within HOMO and LUMO orbitals of LH2NB molecule. The titled molecule has less polar excited state than the ground state as the value of dipole moment is higher in the ground state.

#### 3.4. HOMO-LUMO gap and GCRD analysis

It is very important to study the frontier molecular orbitals (FMOs) as they play a significant character in the reactivity of any molecule. The interactions of HOMO and LUMO in reacting species throughout the progression of chemical reactions are essential among FMOs in stabilization of transition

 Table 1
 The bond lengths [Å], bond angles [°] of LH2NB molecule optimized molecule at HF and B3LYP levels of theory.

Bond lengths (Å)			Bond Angles (°)								
Bonds	Exp. (Natarajan et al., 2012)	HF	B3LYP	Bonds	Exp. (Natarajan et al., 2012)	HF	B3LYP	Bonds	Exp. (Natarajan et al., 2012)	HF	B3LYP
C1015	1.224(2)	1.229	2.256	O15-C1-O16	126.5(2)	132.676	130.927	C23-C24-C26	122.2(2)	121.335	121.318
C1O16	1.266(2)	1.222	1.248	O15-C1-C2	118.0(2)	113.140	113.147	C28-C26-C24	119.7(2)	120.292	120.553
C1–C2	1.530(2)	1.561	1.581	O16-C1-C2	115.5(2)	114.180	115.926	C30-C28-C26	119.7(2)	119.672	119.583
C2-N12	1.483(2)	1.489	1.497	N12-C2-C1	110.7(1)	107.873	105.962	C28-C30-C32	119.8(2)	118.873	118.756
C2–C4	1.533(2)	1.531	1.539	N12-C2-C4	107.8(1)	111.983	111.456	C30-C32-C23	122.2(2)	122.983	123.129
C4–C7	1.485(3)	1.494	1.495	C1C2C4	110.2(2)	113.374	114.156	C30-C32-N33	117.7(2)	115.461	116.628
C7–C8	1.354(3)	1.343	1.370	C7–C4–C2	112.2(2)	110.260	113.151	C23-C32-N33	120.1(2)	121.492	120.175
C7–N14	1.371(3)	1.379	1.388	C8-C7-N14	106.3(2)	106.559	106.153	C10-N13-C8	108.7(2)	109.168	109.883
C8-N13	1.370(3)	1.384	1.383	C8–C7–C4	132.3(2)	132.954	131.590	C10-N14-C7	109.1(2)	109.806	110.163
C10-N13	1.316(3)	1.320	1.341	N14-C7-C4	121.4(2)	120.339	122.256	O36-N33-O37	123.5(2)	124.252	123.343
C10-N14	1.321(3)	1.303	1.332	C7-C8-N13	107.3(2)	106.307	106.596	O36-N33-C32	117.9(2)	117.756	116.947
C22–O34	1.234(3)	1.239	1.252	N13-C10-N14	108.7(2)	108.150	107.202	O37–N33–C32	118.6(2)	117.883	119.630
C22–O35	1.253(2)	1.232	1.265	O34–C22–O35	126.9(2)	127.414	128.482				
C22–C23	1.523(3)	1.522	1.529	O34-C22-C23	116.9(2)	116.333	115.621				
C32-N33	1.459(2)	1.462	1.458	O35-C22-C23	115.9(2)	116.195	115.833				
N33-O37	1.221(3)	1.188	1.225	C24-C23-C22	116.6(2)	119.005	119.570				
N33-O36	1.218(3)	1.199	1.245	C32–C23–C22	127.0(2)	123.946	123.387				

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Bond (H—A) Exp.	. (Natarajan et al. 2012)	HF	B3LYP	CAM-B3LYP	wb97xd
17H—35O 2.18	3	1.7215	1.6015	1.5743	1.6140
20H—34O 1.88	3	1.6902	1.7589	1.7519	1.7490
11H—36O 2.36	5	2.6468	-	_	-
19H—15O 1.97	7	1.9553	1.7709	1.7451	1.7560
18H—15O 1.84	1	2.9384	-	_	-
9H—16O 2.41	1	2.6788	2.0986	2.0722	2.0744

Table 2	Main possible hydrogen	bond in LH2N molecule (	A) obtained different	levels of theory using 6–31G	** basis set.
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structure. Chemical reactivity, kinetic stability, optical polarizability and chemical softness can be described by HOMO– LUMO energy gap values of any molecular system. It is well known that larger the energy gap the harder the material and also shows higher thermal and kinetic stability according to softness–hardness rule. The energy values of HOMO and LUMO orbitals calculated by the B3LYP method are found to be -5.442 and -2.748 eV, respectively. The value of HOMO–LUMO gap is 2.694 eV with chemical hardness 1.347 eV, indicates that the titled material has good chemical stability.

TD-DFT results revealed the excitation energy about 3.778, 3.836 and 3.990 eV calculated by B3LYP, CAM-B3LYP, wb97xd, PBE0/PBE1PBE and MO6 levels of theory viz. corresponding to the electronic transition from ground state to

excited state and displays that the charge transfer from L-Histidine to 2-nitrobenzoate ligand counterparts. Fig. 5 shows the 3-D plot of HOMO–LUMO orbitals and their respective values are presented in Table 6, determined at B3LYP/6-31G\*\* level of theory. The HOMO–LUMO orbitals were also obtained at HF/6-31G\*\* level of theory and their 3-D plot of is shown in Figure 2S (see supplementary data). From the figure it is clearly visible that the HOMO are localized at the lower part of L-Histidine and COO group of 2-nitrobenzoate molecule while LUMO are localized on 2-nitrobenzoate molecule only. LH2NB has an advantage of noteworthy transparency and low absorbance as experimentally proved (Moovendaran et al., 2013).

Further the global chemical reactivity descriptor (GCRD) parameters were calculated as these are used to identify the



Figure 3 Calculated IR and Raman spectra of LH2NB by (a) HF and (b) B3LYP level of theory.

HF Calculated IR freq. (Cm <sup>-1</sup> )	B3LYP Calculated IR freq. (Cm <sup>-1</sup> )	Reported Experimental IR freq. (Cm <sup>-1</sup> ) (Moovendaran et al., 2013)	HF Calculated Raman freq. (Cm <sup>-1</sup> )	B3LYP Calculated Raman freq. (Cm <sup>-1</sup> )	Assignments
3354	3370	-	3355	3369	NH <sub>3</sub> asymmetric stretching
3177	3187	3173	3177	3187	NH symmetric stretching
3105	3153	3138	3113	3153	[NH <sub>3</sub> ] Hydrogen bonded stretching mode
3001,2912,	3075,2980,2962,	2970,2818	3041,3009,2944,2912,2896,2864	3075,2962,	CH <sub>2</sub> , NH <sub>3</sub> symmetric
2864	2928,2893	-		2928,2893	stretching
-	2478	2560,2363	-	2470	C–H combinational overtone
1650,1634	1639	1639	1634	1639	COO <sup>-</sup> asymmetric stretching
1610,1578	1604	1578	1610,1578	1595	NH <sub>3</sub> symmetric stretching
-	1543	1533	1530	1561	N = C - N stretching
1498	1500	1491	1498,1482	-	Ring deformation
1433	1440	1421	1433	1440	COO <sup>-</sup> symmetric stretching
1393	1353	1377	1393,1361	1353	CH <sub>2</sub> deformation
1361,1313 -	1336 	1348	1330	1336 -	NO <sub>2</sub> symmetric stretching + C-C stretching
1265	1301	1288	1265	1310sh,1284	C-C stretching + $C=O$ stretching
1225	1258,1223	1258	1225	1258,1223	C-H in plane bending + C-O bending
1192.1160.1128	1171sh.1137	1136	1144	1172sh.1137	C–H in plane bending
1096	1050	1067	1088	1052sh	Ring breathing, C–H in plane bending
1000,959,935,903	1016,981,938,903	1001,904	1016,959,911	1016,981,938,903	N–H bending
839	834	831	860	834	C–C stretching
815	799	785	823	808	COO <sup>-</sup> bending vibration
750	747,721	754	750	748,721	CCC in plane bending
694	690	694	694	-	NO <sub>2</sub> wagging
662	687	667	662	687	NO <sub>2</sub> rocking
646	653	652	-	-	COO <sup>-</sup> wagging
630	634	629	630	635	Ring deformation
590	600,566	575	574	600sh,566,540	C-NO <sub>2</sub> stretching
509	505	519	509	505	COO <sup>-</sup> wagging
421	419	422	421	419	CCC out of plane bending

**Table 3** HF and B3LYP calculated IR and Raman and experimentally reported IR frequencies with their appropriate assignments for LH2NB using 6-31G<sup>\*\*</sup> basis set.

connection regarding structure, stability and global chemical reactivity of any molecule. GCRD parameters are used in quantitative structure–property, structure–activity, structure–

Table 4	Absorption wavelength ( $\lambda$ ), excitation energies E(eV)
and oscill	ator strengths (f).

Functional	$\lambda$ (nm)	E (eV)	f
B3LYP	328.19	3.778	0.0450
CAM-B3LYP	309.07	4.012	0.0202
wb97xd	310.64	3.991	0.0225
PBE0	317	3.935	0.0377
M06	324	3.820	0.0386

toxicity analysis and a relationship between aromaticity and hardness (Vektariene et al., 2009). DFT has been used in the present study which delivers the interpretations of important universal insights into stability and reactivity of molecular structure (Pearson, 1987). The GCRD (such as hardness ( $\eta$ ), chemical potential ( $\mu$ ), softness (*S*), electronegativity ( $\chi$ ) and electrophilicity index ( $\omega$ )) of LH2NB molecule have been calculated with the help of E<sub>HOMO</sub> as ionization potential (*I*) and E<sub>LUMO</sub> as electron affinity (*I*) respectively. The calculated values of GCRD parameters for LH2NB molecule are presented in Table 6. Thermodynamic properties were also calculated at constant temperature (298.150 K) and pressure (1.000 Atm.) as given in Table 7. The results show that the titled material has good thermal stability.

Table 5 Total dipole moment (Debye) and its components in ground and excited state calculated at different levels of theory.

Components	HF	HF		B3LYP		CAM-B3LYP		wb97xd	
	G.S.	E.S.	G.S.	E.S.	G.S.	E.S.	G.S.	E.S.	
$\mu_X$	-7.434	-5.732	-8.761	-3.292	-8.586	-3.329	-8.582	-3.441	
$\mu_{v}$	14.268	6.389	12.947	11.167	13.384	11.058	13.358	11.553	
$\mu_z$	-9.111	1.413	-8.358	-1.806	-8.447	-2.053	-8.492	-2.076	
$\mu_{tot}$	18.489	8.699	17.727	11.782	18.006	11.730	18.005	12.232	



Figure 4 Calculated optical absorption spectra of LH2NB molecule at different levels of theory.

## 3.5. *NLO* (polarizability and first order hyperpolarizability) effect

The NLO processes are of technological significance in the production of laser devices, therefore in recent decade it is considered as a key subject of study, in which hyperpolarizability played a vital role and maintains the growing application and adoring to the exact determination of this property (Bishop, 1994; Prasad and Williams, 1991). As per the available literature numerous current studies display that besides a suitable action of electron correlation and a careful selection of the basis set function, the calculations of molecular electrical properties like polarizability and hyperpolarizability the inclusion of contributions arising from the nuclear motion is of fundamental importance (Ingamells et al., 1998; Raptis et al., 1999; Eckart et al., 2001).

Author know that the importance of the polarizability and first hyperpolarizability values of any molecule is reliant on the efficiency of electronic communication between acceptor and donor groups as these are the key elements in molecular charge transfer (Udayakumar et al., 2011; Wolinski et al., 1990; Cheeseman et al., 1996; Kalinowski et al., 1988; Pihlaja et al., 1994). To calculate the linear and nonlinear optical properties of LH2NB molecule author have applied B3LYP, CAM-B3LYP and wb97xd levels of theory using 6-31G\*\* basis set. The calculated values of electronic dipole moment ( $\mu$ ), total average polarizability ( $\alpha_{tot}$ ), anisotropy of polarizability ( $\Delta \alpha$ ), static first hyperpolarizability  $(\beta_0)$  and total first hyperpolarizability  $(\beta_{tot})$  are presented in Table 7 at B3LYP/6-31G<sup>\*\*</sup> level of theory. The above said values calculated by other methods are presented in Tables 8-10.



Figure 5 The 3-D plot of HOMO–LUMO orbitals of LH2NB molecule at isoval =  $\pm 0.02$  a.u.



**Figure 6** Variation of  $\beta_{tot}$  values calculated at different methods in comparison with urea.

The different mathematical expressions are used to calculate the above parameters as expressed below:

Total dipole moment is calculated by:

$$\mu_{tot} = \left(\mu_x^2 + \mu_y^2 + \mu_z^2\right)^{\frac{1}{2}}$$
(9)

Total polarizability is calculated by:

$$\alpha_{tot} = \frac{1}{3} (\alpha_{xx} + \alpha_{yy} + \alpha_{zz}) \tag{10}$$

Table 6	The calc	ulated en	nergy	values	of	frontier	molecular
orbitals	(FMOs),	energy	differ	ence a	and	global	reactivity
descripto	ors at B3L	YP/6-310	3** lev	vel of the	heo	ry.	

	B3LYP	
Orbitals	a.u.	eV
E <sub>HOMO</sub>	-0.200	-5.442
E <sub>HOMO-5</sub>	-0.243	-6.612
E <sub>HOMO-1</sub>	-0.211	-5.742
E <sub>LUMO</sub>	-0.101	-2.748
E <sub>LUMO+1</sub>	-0.058	-1.578
$\Delta E_{HOMO}$ -LUMO	0.099	2.694
$\Delta E_{HOMO-5}$ -LUMO	0.142	3.864
$\Delta E_{HOMO-1} - LUMO + 1$	0.153	4.163
η	0.05	1.347
μ	-0.151	-4.095
S	0.074	2.020
χ	0.151	4.095
	0.229	6.225

Anisotropy of polarizability is calculated by:

$$\Delta \alpha = \frac{1}{\sqrt{2}} \times \sqrt{\left[ \left( \alpha_{xx} - \alpha_{yy} \right)^2 + \left( \alpha_{yy} - \alpha_{zz} \right)^2 + \left( \alpha_{zz} - \alpha_{xx} \right)^2 + 6\alpha_{xz}^2 \right]}$$
(11)

The components of first hyperpolarizability can be calculated using the following relation:

$$\beta_i = \beta_{iii} + \sum_{i \neq j} \left[ \frac{(\beta_{ijj} + 2\beta_{jii})}{3} \right]$$
(12)

Using x, y, z components the magnitude of first hyperpolarizability ( $\beta_{tot}$ ) can be achieved by:

$$\beta_{tot} = \sqrt{(\beta_x^2 + \beta_y^2 + \beta_z^2)} \tag{13}$$

where 
$$\beta_x$$
,  $\beta_y$  and  $\beta_z$  are:

$$p_x = (p_{xxx} + p_{xxy} + p_{xyy})$$
$$\beta_y = (\beta_{yyy} + \beta_{xxz} + \beta_{yyz})$$
$$\beta_z = (\beta_{xzz} + \beta_{yzz} + \beta_{zzz})$$

So, the final equation for magnitude of total first hyperpolarizability calculation is given by:

$$\beta_{tot} = \sqrt{\left[ (\beta_{xxx} + \beta_{xxy} + \beta_{xyy})^2 + (\beta_{yyy} + \beta_{xxz} + \beta_{yyz})^2 + (\beta_{xzz} + \beta_{yzz} + \beta_{zzz})^2 \right]}$$
(14)

Table 7	Thermodynamic parameters of LH2NB molecule.						
Methods	E (Thermal) (KCal/Mol)	C <sub>V</sub> (Cal/Mol- Kelvin)	S (Cal/Mol- Kelvin)				
HF	204.003	71.220	151.059				
B3LYP	189.505	76.461	153.956				
CAM-	192.504	75.144	152.836				
B3LYP							
Wb97xd	192.562	74.984	149.904				

**Table 8** The calculated values of polarizability ( $\alpha$ ), dipole moment ( $\mu$ ) and hyperpolarizability ( $\beta$ ) along their individual tensor components of LH2NB at B3LYP/6-31G<sup>\*\*</sup> level of theory.

Components	a.u.	$esu (\times 10^{-24})$	Component	a.u.	esu (×10 <sup>-30</sup> )
$\alpha_{xx}$	247.415	36.667	$\beta_{xxx}$	26.530	0.229
$\alpha_{xy}$	-4.757	-0.705	$\beta_{xxy}$	458.966	3.960
$\alpha_{\nu\nu}$	205.029	30.385	$\beta_{xvv}$	29.511	0.255
$\alpha_{xz}$	3.504	0.519	$\beta_{vvv}$	706.008	6.092
$\alpha_{yz}$	-6.248	-0.926	$\beta_{xxz}$	-39.513	-0.341
$\alpha_{zz}$	97.932	14.514	$\beta_{xvz}$	35.165	0.303
$\alpha_{tot}$	183.459	27.189	$\beta_{yyz}$	24.906	0.215
Δα	133.576	19.795	$\beta_{xzz}$	22.689	0.196
$\mu_x$	1.296	-3.292D	$\beta_{vzz}$	22.695	0.196
$\mu_v$	4.395	11.163D	$\beta_{zzz}$	-6.459	-0.056
$\mu_z$	0.711	-1.805D	$\beta_0$	517.805	4.468
$\mu_{tot}$	4.637	11.777D	$\beta_{tot}$	863.008	7.447
$\mu_{urea}$	0.541	1.3732D (Govindarasu and Kavitha, 2014;	$\beta_{urea}$	43.203	0.3728 (Govindarasu and Kavitha,
		Shkir and et al., 2015b)			2014; Shkir and et al., 2015b)

**Table 9** The calculated values of polarizability ( $\alpha$ ), hyperpolarizability ( $\beta$ ) and dipole moment ( $\mu$ ) along their individual tensor components of LH2NB calculated at CAM-B3LYP/6-31G<sup>\*\*</sup> level of theory.

Components	a.u.	esu (×10 <sup>-24</sup> )	Component	a.u.	esu (×10 <sup>-30</sup> )
$\alpha_{xx}$	230.762	34.199	$\beta_{xxx}$	-140.742	-1.215
$\alpha_{xy}$	-2.754	-0.408	$\beta_{xxy}$	211.661	1.826
$\alpha_{vv}$	192.930	28.592	$\beta_{xvv}$	3.060	0.026
$\alpha_{xz}$	3.726	0.552	$\beta_{yyy}$	421.943	3.64
$\alpha_{vz}$	-7.240	-1.073	$\beta_{xxz}$	-40.938	-0.353
azz	100.936	14.958	$\beta_{xyz}$	28.431	0.245
$\alpha_{tot}$	174.876	25.916	$\beta_{vvz}$	26.563	0.229
Δα	115.993	17.19	$\beta_{xzz}$	16.941	0.146
$\mu_x$	1.404	-3.329D	$\beta_{vzz}$	19.055	0.164
$\mu_v$	4.462	11.058D	$\beta_{zzz}$	-8.906	-0.077
$\mu_z$	0.820	-2.053D	$\beta_0$	249.067	2.149
$\mu_{tot}$	4.749	11.730D	$\beta_{tot}$	415.112	3.582
$\mu_{urea}$	0.541	1.3732D (Govindarasu and Kavitha,	$\beta_{urea}$	43.203	0.3728 (Govindarasu and Kavitha,
		2014; Shkir and et al., 2015b)			2014; Shkir and et al., 2015b)

First hyperpolarizability can be labeled as a  $3 \times 3 \times 3$ matrix as it is known as a third rank tensor and 27 components of the 3D matrix are reduced to 10 components due to the Kleinman symmetry (Kleinman, 1962), and can be specified in lower tetrahedral format as it is evident that the lower part of  $3 \times 3 \times 3$  matrices is tetrahedral. Further, the static first hyperpolarizability ( $\beta_0$ ) was also calculated using the following relation  $\beta_0 = \frac{3}{5}\beta_{tot}$  and given in Table 7. The electronic communication of two different parts of any system plays a key role in polarizability and hyperpolarizability values. Calculated values of all components and their resultants of dipole moment, polarizability and hyperpolarizability by B3LYP/ 6-31G<sup>\*\*</sup> are given in Table 7 (in a.u and esu (for  $\alpha$ , 1 a.u. =  $0.1482 \times 10^{-24}$  esu and for  $\beta$ , 1 a.u. =  $0.008629 \times 10^{-30}$  esu)) (Govindarasu and Kavitha, 2014; Shkir et al., 2015b). The value of  $\mu_{tot}$  and  $\beta_{tot}$  at molecular level are found to be 8 times and 20 times higher than urea ( $\mu_{Urea} = 1.3732D$  and  $\beta_{urea} = 0.3728 \times 10^{-30}$  esu) (Govindarasu and Kavitha, 2014; Shkir et al., 2015b) respectively. These values are higher than many other organic, semiorganic and inorganic materials also (Shkir et al., 2014, 2015b; Adant et al., 1995; Karabacak and Cinar, 2012; Raju et al., 2015; Abbas et al. 2015; Singh et al., 2012; Govindarasu et al., 2014). Further, these values were calculated by exchange correlation functional CAM-B3LYP and wb97xd level of theory with 6-31G\*\* basis set and presented in Table 8 and 9. In recent reports the experimental value of second harmonic generation of LH2NB was found to be 2 times higher than KDP at bulk level (Moovendaran et al., 2013; Natarajan et al., 2012). The representation of variation of first order hyperpolarizability calculated at all the methods with reference to urea is shown in Fig. 6. Therefore, from both experimental and theoretical studies it is evident that the titled compound is a very good nonlinear optical material and can be useful for laser device applications.

#### 3.6. Molecular electrostatic potential (MEP) analysis

Fig. 7 shows the MEP (3-D plot) of LH2NB molecule. MEP plays a very important role in understanding many points such as: (i). About its positive and negative regions as nucleophile and electrophile at molecular level(Murray and Sen, 1996; Scrocco and Tomasi, 1978), (ii). It is a measurement of electrostatic potential on constant electron density surface, (iii). It is

Components	a.u.	esu (×10 <sup>-24</sup> )	Component	a.u.	esu (×10 <sup>-30</sup> )
$\alpha_{xx}$	233.588	34.618	$\beta_{xxx}$	-144.983	-1.251
$\alpha_{xv}$	-3.531	-0.523	$\beta_{xxy}$	243.199	2.099
$\alpha_{vv}$	194.649	28.847	$\beta_{xvv}$	-10.829	-0.093
a <sub>xz</sub>	3.173	0.47	$\beta_{vvv}$	413.895	3.571
$\alpha_{_{VZ}}$	-7.131	-1.057	$\beta_{xxz}$	-35.636	-0.307
azz	99.735	14.781	$\beta_{xvz}$	26.934	0.232
$\alpha_{tot}$	175.991	26.082	$\beta_{yyz}$	31.685	0.273
Δα	119.377	17.692	$\beta_{xzz}$	20.229	0.175
$\mu_x$	1.354	-3.441D	$\beta_{vzz}$	18.165	0.156
$\mu_{v}$	4.545	11.553D	$\beta_{zzz}$	-9.987	-0.086
$\mu_z$	0.817	-2.076D	$\beta_0$	252.070	2.175
$\mu_{tot}$	4.812	12.232D	$\beta_{tot}$	420.116	3.625
$\mu_{urea}$	0.541	1.3732D (Govindarasu and Kavitha,	$\beta_{urea}$	43.203	0.3728 (Govindarasu and Kavitha,
		2014; Shkir and et al., 2015b)			2014; Shkir and et al., 2015b)

**Table 10** The calculated values of polarizability ( $\alpha$ ), hyperpolarizability ( $\beta$ ) and dipole moment ( $\mu$ ) along their individual tensor components of LH2NB calculated at wb97xd/6-31G<sup>\*\*</sup> level of theory.



Figure 7 MEP plot of LH2NB molecule at ISO value 0.02.

also a very useful property to examine the reactivity of molecular species by guessing that either the approaching nucleophile is attracted to a positive region of the molecular system (iv). It will also help to know the simultaneous information about the molecular size, shape along with its positive, negative and neutral electrostatic potential regions in terms of color grading (Murray and Sen, 1996; Lowdin, 1979). The MEP plot was also determined at HF/6-31G<sup>\*\*</sup> level of theory and shown in Figure 3S (see supplementary data). The red and blue colors on surface shows highest negative and positive potential respectively.

The COO group of L-Histidine shows the negative potential as represented by red as well as yellow colors which is favorable for electrophilic attack while near the other atoms, it has positive potential with blue color and favorable for nucleophilic attack. The yellow color on the surface shows intermediate negative potential. The green color represents the slight positive or neutral potential over the carbon atoms.

#### 4. Conclusions

The stable geometry of LH2NB molecule has been achieved and the optimized structural parameters were found in good agreement with the reported experimental values. The calculated vibrational modes of LH2NB are found in strong agreement with earlier experimental reports. The excitation energy of LH2NB molecule was calculated by TD-B3LYP, TD-CAM-B3LYP, TD-wb97xd, TD-PBE0 and TD-MO6 using 6-31G\*\* basis set, and found to be 3.778, 4.012, 3.991, 3.935 and 3.820 eV and with an oscillatory strength of 0.045, 0.020, 0.023, 0.0377 and 0.0386 respectively. The charge interaction taking place within the molecule has been studied from HOMO-LUMO orbitals and energy gap was calculated. The chemical and thermal stability of LH2 NB has been studied with help of global chemical reactivity descriptors and thermodynamic parameters. The value of dipole moment is found to be lower in the excited state than ground state and shows that the ground state is more polar than excited state. The  $\mu_{tot}$  and  $\beta_{tot}$  values are found to be 8 and 20 times higher than standard urea molecule at molecular level respectively. The molecular electrostatic potential plot shows that the negative potential sites are on electronegative oxygen atoms of COO group and positive potential sites are on hydrogen atoms. The higher total first hyperpolarizability indicates that the studied compound is maybe an excellent candidate for laser device applications.

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#### Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at http://dx.doi.org/10.1016/j.jksus.2016. 03.002.

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