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An efficient modification of the decomposition method with a convergence parameter for solving Korteweg de Vries equations

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ABSTRACT

In the present paper, an efficient modification the convergence parameter based on the Adomian Decomposition Method (ADM) is proposed and investigated for a class of nonlinear evolution equations; specifically, the Korteweg de Vries (KdV) equations. We show that the proposed analysis possesses increased accuracy when compared to the standard ADM. Moreover, the optimal value of such a convergence parameter is determined by minimizing the averaged residual error. For such a convergence parameter value, an approximate solution is found to be closer to the available exact solution than the corresponding approximate solution without a convergence parameter for the same number of solution components. The approach proposed may be readily extended to other nonlinear differential and integral equations.

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1. Introduction

The well-known evolution equations that occur frequently in mathematical physics and shallow water applications can be formulated as

$$u_t = \sum_{m=0}^M a_m \frac{\partial^m u}{\partial x^m} + \sum_{m=0}^N b_m \frac{\partial^m (u^{k+1})}{\partial x^m} + C \frac{\partial^{i+1} u}{\partial x^i \partial t} + f(x, t), \quad (1)$$

where $f(x, t)$ is the homogeneous term function; C, a_m, b_m real constants; k and i natural numbers and M and N are nonnegative integers. The exact solution of Eq. (1) is always very difficult to find (Yusuf et al., 2018) which necessitates the study of its special cases including the Korteweg de Vries (KdV) equation, the Benjamin–Bona–Mahony equation (BBM), the Burger's equation and the regularized long-wave equation (RLWE) among others. The well-known KdV equation reads

$$u_t + ruu_x + u_{3x} = f(x, t), \quad x \in [a, b], \quad t \geq 0, \quad (2)$$

where r is a constant mostly preferred 1 or 6 for solution description, see (Soliman and Abdou, 2008; Syam, 2005; Varley and Seymour, 1998; Wazwaz, 1999; Wazwaz, 2001). Recently, the Adomian decomposition method (ADM) (Adomian, 1983; Adomian, 1986; Adomian, 1988; Adomian, 1991; Adomian and Rach, 1992) has been broadly used in treating a variety of mathematical problems mostly modeled in nonlinear differential and integral equations (Aly et al., 2012; AlQarni et al., 2016; Bakodah, 2012; Bakodah, 2012; Bakodah et al., 2017; Bakodah et al., 2015; Bakodah and Darwish, 2013; Bakodah et al., 2016; Banaja et al., 2017; Bulut et al., 2013); see also (Sabi'u et al., 2018; Ebaid, 2011; Ebaid et al., 2015; Duan and Rach, 2011; Duan et al., 2012; Duan, 2010; Inc et al., 2018; Nuruddeen et al., 2018; Qureshi and Ramos, 2018) for various modifications of the ADM and (Liao, 2010; Rach, 1987; Schiesser, 1994; Abdel-Gawad et al., 2018; Aliya et al., 2018; Ansari et al., 2018; Yusuf et al., 2018; Zhang and Liang, 2016) for other methods for solving various evolution equations, respectively. The main part of the ADM is generating the Adomian's polynomials

$$A_n = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} N \left(\sum_{i=0}^n \lambda^i u_i \right) \right]_{\lambda=0}, \quad n \geq 0, \quad (3)$$

with A_n denoting the Adomian's polynomials of degree n for the nonlinear term appearing in the considered equation while

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$u = \sum_{i=0}^{\infty} u_i(x, t)$ is the series solution converging to the exact one. Furthermore, in the case of the Adomian's polynomials for multi-variable nonlinearities and differential nonlinearities, we have

$$A_n = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} N \left(\sum_{i=0}^n \lambda^i u_{1,i}, \dots, \sum_{i=0}^n \lambda^i u_{m,i} \right) \right]_{\lambda=0}, \quad n \geq 0,$$

where $N(u_1, \dots, u_m)$ is the multivariable nonlinearity of m arguments. Again, both the one and multivariable Adomian's polynomials can respectively be rapidly generated to high orders by algorithms and Mathematica subroutines crafted by Duan (Inc et al., 2018). Moreover, Zhang and Liang (2016) recently proposed a new convergence parameter for the ADM that controls the convergence-region and the rate of the optimal series solution. However, the recent modification of the ADM based on Zhang and Liang (2016) will be proposed in this paper for a class of evolution equations, particularly the KdV equations. Some special test problems of interest would be numerically analyzed by the proposed scheme and provide the error estimate and analysis. The extension of the domain of convergence which provides more accurate approximations can be regarded as the main advantage of the scheme.

2. Analysis of the standard ADM

Considering the linear operators

$$L_{x^m} = \frac{\partial^m}{\partial x^m} \text{ and } L_t = \frac{\partial}{\partial t};$$

Eq. (1) can be expressed as

$$L_t u = \sum_{m=0}^M a_m L_{x^m} u + \sum_{m=0}^N b_m L_{x^m} u^{k+1} + C L_t L_{x^i} u + f(x, t), \quad (4)$$

with $L_t^{-1}(.) = \int_0^t (.) dt$ as the inversion operator of L_t .

Applying the above inversion operator coupled to initial data $u(x, 0) = h(x)$ on Eq. (4), we obtain

$$u = h(x) + L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} u + \sum_{m=0}^N b_m L_{x^m} u^{k+1} + C L_t L_{x^i} u + f(x, t) \right). \quad (5)$$

Thus, the standard ADM offers the following sequential recursion solution of Eq. (1) by decomposing $u(x, t)$ to infinite series $\sum_{n=0}^{\infty} u_n(x, t)$ as

$$u_0(x, t) = h(x) + L_t^{-1}(f(x, t)),$$

$$u_{n+1}(x, t) = L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} u_n + \sum_{m=0}^N b_m L_{x^m} A_n + C L_t L_{x^i} u_n \right), \quad n \geq 0, \quad (6)$$

where $A_n, n \geq 0$, are the Adomian polynomials (Adomian, 1983; Adomian, 1986; Adomian, 1988; Adomian, 1991).

3. Analysis of the ADM with a convergence parameter

In this section, the concept of a convergence parameter Zhang and Liang (2016) shall be applied to derive the general recursive scheme of Eq. (1) and later demonstrate its application to solve several special cases in the next section. To begin with, a convergence-control parameter c and an artificial parameter ϵ are introduced to Eq. (4) which reads

$$L_t u = (\epsilon c + \epsilon^2(1 - c)) \left(\sum_{m=0}^M a_m L_{x^m} u + \sum_{m=0}^N b_m L_{x^m} u^{k+1} + C L_t L_{x^i} u \right) + f(x, t). \quad (7)$$

One then set

$$u(x, t) = \sum_{i=0}^{\infty} v_i(x, t, c) \epsilon^i. \quad (8)$$

Applying the inversion operator L_t^{-1} on Eq. (7), we get

$$u(x, t, c, \epsilon) = (\epsilon c + \epsilon^2(1 - c)) L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} u + \sum_{m=0}^N b_m L_{x^m} u^{k+1} + C L_t L_{x^i} u \right) + h(x) + L_t^{-1} f(x, t). \quad (9)$$

Now, putting Eq. (8) into Eq. (9) and bringing together the coefficients of like-powers of ϵ , we obtain the following recursive scheme:

$$\begin{aligned} v_0(x, t, c) &= h(x) + L_t^{-1} f(x, t), \\ &= c L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} v_0 + \sum_{m=0}^N b_m L_{x^m} B_0 + C L_t L_{x^i} v_0 \right), \\ &= c L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} v_{n-1} + \sum_{m=0}^N b_m L_{x^m} B_{n-1} + C L_t L_{x^i} v_{n-1} \right) \\ &\quad + (1 - c) L_t^{-1} \left(\sum_{m=0}^M a_m L_{x^m} v_{n-2} + \sum_{m=0}^N b_m L_{x^m} B_{n-2} + C L_t L_{x^i} v_{n-2} \right), \\ &n \geq 2, \end{aligned} \quad (10)$$

with B_n the Adomian's polynomials to be computed using

$$B_n = \frac{1}{n!} \left[\frac{d^n}{d\epsilon^n} N \left(\sum_{i=0}^n \epsilon^i v_i(x, t, c) \right) \right]_{\epsilon=0}, \quad n \geq 0.$$

When $\epsilon = 1$ in Eq. (8), the solution $u(x, t)$ with convergence-parameter c is presented by

$$u(x, t) = \sum_{i=0}^{\infty} v_i(x, t, c). \quad (11)$$

Then we calculate an n^{th} -order approximation to the solution as

$$\varphi_{m+1}(x, t, c) = \sum_{n=0}^m v_n(x, t, c).$$

To extend the domain of convergence or rather obtain more accurate approximations; the value of the parameter c is to be determined using the discrete averaged square residual error defined as (Nuruddeen et al., 2018)

$$E_m(x, t, c) = \frac{1}{k} \sum_{i=1}^k \frac{1}{k} \sum_{j=1}^k \left[F \left(\sum_{n=0}^m v_n(j\Delta x, i\Delta t, c) \right) \right]^2. \quad (12)$$

Further, the averaged residual error E_m against the determined value of c gives the optimal value of c when plotted. We will demonstrate later that the new approach is easily implemented on any symbolic software such as Maple or Mathematica. Also, the optimal value of c is obtained via solving the differential equation

$$\frac{\partial E(x, t, c)}{\partial c} = 0.$$

Finally, it can be remarked that the optimal choice of c could significantly improve the convergence region and rate of the series solution. We show the effectiveness of the present approach by investigating certain test KdV and coupled KdV systems.

4. Applications

The present section demonstrates the high-accuracy of presented method by considering certain KdV and coupled KdV sys-

tems of type Eq. (1). We will present the numerical results based on our method and demonstrate its accuracy by studying the absolute error and the number of iterations involved.

Example 1. Consider the KdV equation with the given initial data:

$$u_t - 6uu_x + u_{3x} = 0, |t| < 1, \quad (13a)$$

$$u(x, 0) = \frac{1}{6}(x - 1). \quad (13b)$$

Here, $v_0(x, 0, c)$ is selected as $v_0(x, t, c) = g(x, t) = \frac{1}{6}(x - 1)$, and consequently Eq. (10) leads to

$$v_1(x, t, c) = c \int_0^t [(v_0(x, t, c))_{3x} - aB_0] dt,$$

$$v_2(x, t, c) = c \int_0^t [(v_1(x, t, c))_{3x} - aB_1] dt + (1 - c) \int_0^t [(v_0(x, t, c))_{3x} - aB_0] dt.$$

Using the parameterized recursion scheme in Eq. (10) with the appropriate Adomian polynomials yields the first several solution components as

$$v_1(x, t) = \frac{1}{6}ct(-1 + x), \quad (14)$$

$$v_2(x, t) = \frac{1}{6}t(1 - c + c^2t)(-1 + x), \quad (15)$$

$$v_3(x, t) = \frac{1}{6}ct^2(2 - 2c + c^2t)(-1 + x), \quad (16)$$

$$v_4(x, t) = \frac{1}{6}t^2(1 - 2c - 3c^2t + c^4t^2 + c^2(1 + 3t))(-1 + x), \quad (17)$$

and so on. Substituting Eqs. (14)–(17) into Eq. (11) gives the solution $u(x, t)$, from which it can easily be verified that the closed form solution when $c = 0$ is in this form:

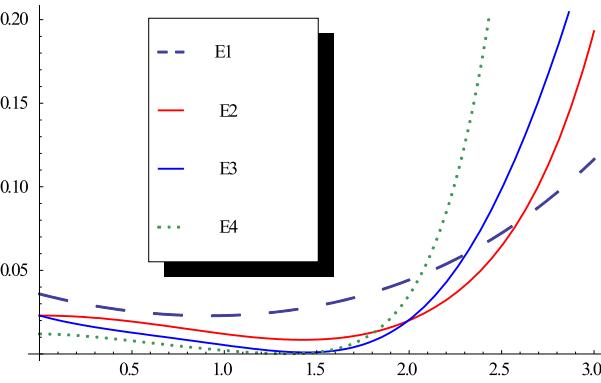


Fig. 1. The residual errors at $m = 1, 2, 3, 4$, and 5.

$$u(x, t) = \frac{1}{6} \frac{(x - 1)}{(1 - t)}, |t| < 1; \quad (18)$$

Also, to validate the obtained solutions, the curves of the discrete averaged square residual errors E_m versus c are shown in Fig. 1. This figure points out that the optimal value of c is about 1.3392966.

The results produced by the proposed method are overall more accurate than the standard ADM, as shown in Table 1a.

The results in Table 1a reveal that the proposed approach is more accurate than the standard ADM in most cases for the chosen values for x and t (Fig. 2).

Remark. As shown in Table 1a, the numerical results of the proposed method are closer to the values obtained from the analytical solution. Moreover, as shown in Table 1b that the error turned out to be smaller as the number of iterations increased.

Example 2. To further illustrate the effectiveness of the proposed method, two versions of the modified KdV (mKdV) equations are considered:

(i) The first version is given by the following equation

$$u_t + 6u^2u_x + u_{3x} = 0, \quad (19)$$

with the initial condition

$$u(x, 0) = a - \frac{4a}{4a^2x^2 + 1}, \quad (20)$$

where a is any real constant. On applying the analysis of the previous example, the first several solution components are computed as

$$v_0(x, t) = a - \frac{4a}{1 + 4a^2x^2},$$

$$v_1(x, t) = -\frac{192a^5ctx}{(1 + 4a^2x^2)^2},$$

$$v_2(x, t) = -\frac{192a^5t(x - cx + 36a^4c^2tx^2 + a^2(-3c^2t - 4(-1 + c)x^3))}{(1 + 4a^2x^2)^3},$$

$$v_3(x, t) = -\frac{1}{(1 + 4a^2x^2)^4} 1152a^7ct^2(-1 + 8a^2x^2 + 48a^4x^4 + 48a^4c^2tx(-1 + 4a^2x^2) + c(1 - 8a^2x^2 - 48a^4x^4)),$$

and so on. The other calculated terms give the closed form solution when $c = 0$ as

$$u(x, t) = a - \frac{4a}{(4a^2(x - 6a^2t)^2 + 1)}.$$

Fig. 3 gives the plots of the averaged residual errors E_m against c showing the optimal value of c around 1.001.

Table 1a

Absolute errors of the proposed method compared with absolute errors of the classical ADM, where $t \in \{0, 0.3, 0.5\}$ and $x \in [0.1, 0.5]$, $c = 1.3392966$.

x/t	Present		Standard ADM		Present		Standard ADM	
	0.2	0.5	0.2	0.5	0.6	0.6	0.6	0.6
0.1	0.0000360	0.0000600	0.0034728	0.0093750	0.0002255	0.029160		
0.2	0.0000320	0.0000533	0.0030869	0.0083333	0.0002004	0.025920		
0.3	0.0000280	0.0000466	0.0027011	0.0072916	0.0001754	0.022679		
0.4	0.0000240	0.0000400	0.0023152	0.0062500	0.0001503	0.019440		
0.5	0.0000200	0.0000333	0.0019293	0.0052083	0.0001253	0.016199		

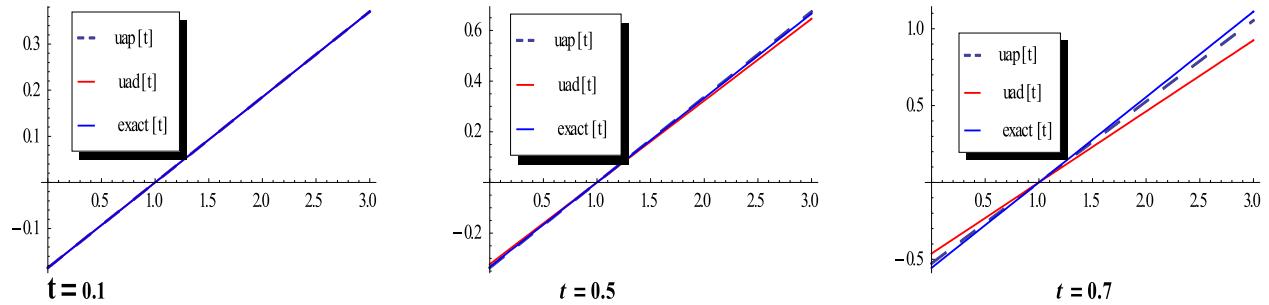


Fig. 2. The plots of the exact, proposed method and the ADM solutions, respectively at $t = 0.1, 0.5$ and 0.7 .

Table 1b

The absolute errors of the proposed method against the number of components at $x = 0.5$.

n	$t = 0.1$	$t = 0.4$
1	0.00260878	0.00808335
2	0.00076428	0.00482109
3	0.00001800	0.00411788
4	0.00003267	0.00152784

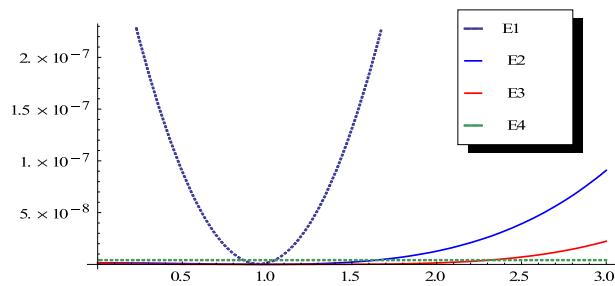


Fig. 3. The residual errors at $m = 1, 2, 3$, and 4 .

The results produced by our method are put side by side with those obtained by the standard ADM and listed in **Table 2**. The profile of the solitary wave at $t = 0.5$ is displayed in **Fig. 4**.

(ii) The second version of the mKdV equation has the traveling wave solution with the initial data:

$$u(x, 0) = \sqrt{a} \operatorname{sech}(k + \sqrt{a}x), \quad (21)$$

with k constant and for $a \geq 0$.

Repeating the previous analysis, we have

$$v_0(x, t) = \sqrt{a} \operatorname{Sech}(k + \sqrt{a}x),$$

$$v_1(x, t) = a^2 ct \operatorname{Sech}(k + \sqrt{a}x) \operatorname{Tanh}(k + \sqrt{a}x),$$

$$v_2(x, t) = \frac{1}{4} a^2 t \operatorname{Sech}(k + \sqrt{a}x)^3 (a^{3/2} c^2 t (-3 + \operatorname{Cosh}[2(k + \sqrt{a}x)]) - 2(-1 + c) \operatorname{Sinh}[2(k + \sqrt{a}x)]),$$

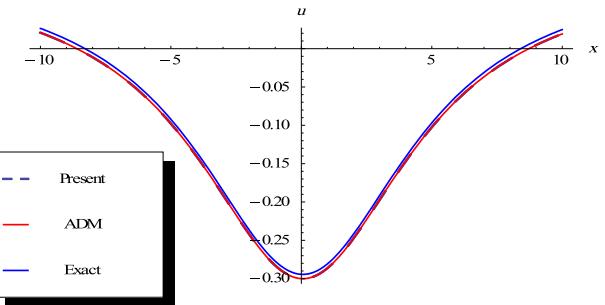


Fig. 4. The plots of the exact, proposed method and the ADM solutions, respectively at $t = 1.0$.

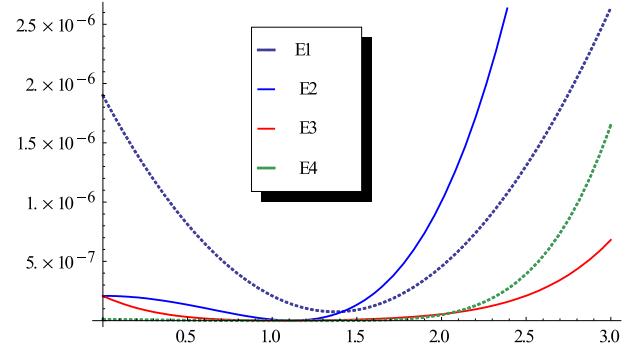


Fig. 5. The residual errors at $m = 1, 2, 3$, and 4 .

$$v_3(x, t) = \frac{1}{24} a^{7/2} c t^2 \operatorname{Sech}(k + \sqrt{a}x)^4 (30(-1 + c) \operatorname{Cosh}(k + \sqrt{a}x) - 6(-1 + c) \operatorname{Cosh}[3(k + \sqrt{a}x)] + a^{3/2} c^2 t (-23 \operatorname{Sinh}[k + \sqrt{a}x] + \operatorname{Sinh}[3(k + \sqrt{a}x)])),$$

and so on. At $c = 0$, the other calculated terms give the following exact solution:

$$u(x, t) = \sqrt{a} \operatorname{Sech}[k + \sqrt{a}(x - at)]. \quad (22)$$

Table 2

Absolute errors of the proposed method compared with absolute errors of the classic ADM, where $t \in \{0.3, 0.5\}$ and $x \in [0, 1, 0.5]$, $c = 1.0003$.

x/t	Present		Standard ADM		Present		Standard ADM	
	0.3	0.5	0.3	0.5	0.3	0.5	0.3	0.5
0.1	1.554×10^{-15}		9.436×10^{-16}		2.665×10^{-15}		1.854×10^{-14}	
0.2	1.998×10^{-15}		7.771×10^{-16}		7.216×10^{-15}		1.781×10^{-14}	
0.3	2.276×10^{-15}		7.216×10^{-16}		1.171×10^{-15}		1.676×10^{-14}	
0.4	2.498×10^{-15}		7.216×10^{-16}		1.593×10^{-15}		1.559×10^{-14}	
0.5	2.831×10^{-15}		4.996×10^{-16}		2.015×10^{-15}		1.382×10^{-14}	

Table 3

Absolute errors of the present method compared with absolute errors of the classical ADM, where $t \in \{0.3, 0.5\}$ and $x \in [0.1, 0.5]$, $c = 1.0273$.

x/t	Present	Standard ADM	Present	Standard ADM
	0.3		0.5	
0.1	6.418×10^{-10}	1.744×10^{-9}	4.241×10^{-9}	3.851×10^{-8}
0.2	5.807×10^{-10}	1.578×10^{-9}	3.835×10^{-9}	3.485×10^{-8}
0.3	5.254×10^{-10}	1.428×10^{-9}	3.469×10^{-9}	3.154×10^{-8}
0.4	4.754×10^{-10}	1.293×10^{-9}	3.137×10^{-9}	2.854×10^{-8}
0.5	4.301×10^{-10}	1.170×10^{-9}	2.838×10^{-9}	2.582×10^{-8}

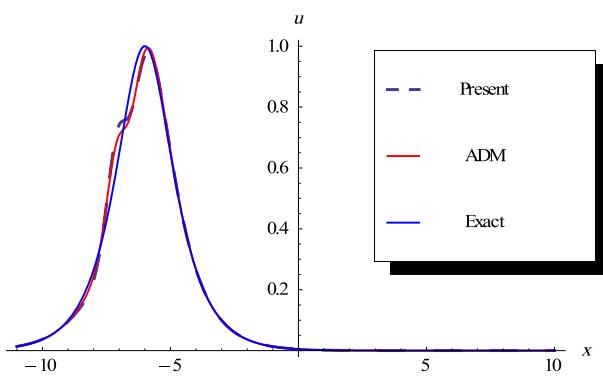


Fig. 6. The plots of the exact, proposed method and the ADM solutions, respectively at $t = 1.0$.

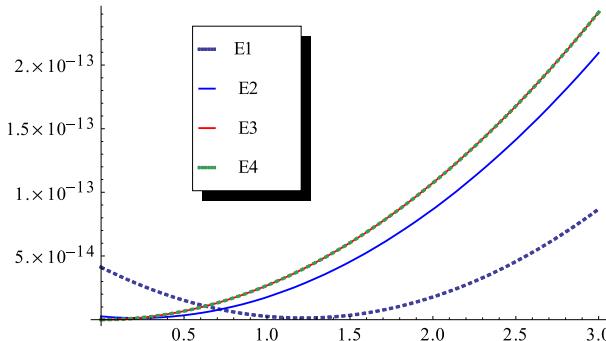


Fig. 7. The residual errors at $m = 1, 2, 3$, and 4.

Here also, Fig. 5 gives the plots of the averaged residual errors E_m against c showing the optimal value of c around 1.03.

The results derived by the present method are much better than those obtained by the standard ADM as showed in Table 3. In addition, the profile of the solitary wave at time $t = 0.5$ is compared in Fig. 6.

Example 3. Consider the combined KdV-mKdV equation

$$u_t + r_1 u u_x + r_2 u^2 u_x + u_{3x} = 0, \quad (23)$$

Table 4

Absolute errors of the proposed method compared with absolute errors of the classical ADM, where $t \in \{0.9, 1.0\}$ and $x \in [0.1, 0.5]$, $c = 0.556$.

x/t	Present	Standard ADM	Present	Standard ADM
	0.9		1.0	
0.1	6.6613×10^{-7}	0.00000104	3.832×10^{-6}	0.00000127
0.2	5.9739×10^{-7}	0.00000115	9.763×10^{-7}	0.00000139
0.3	5.2864×10^{-7}	0.00000126	8.904×10^{-7}	0.00000152
0.4	6.207010^{-7}	0.00000137	8.045×10^{-7}	0.00000164
0.5	5.4336×10^{-7}	0.00000148	7.185×10^{-7}	0.00000176

where r_1 and r_2 are constants with the following initial data:

$$u(x, 0) = \alpha + \gamma \operatorname{csch}(kx), k = \sqrt{\beta}, \alpha = -\frac{r_1}{2r_2} \text{ and } \gamma = \sqrt{\frac{6\beta}{r_2}}.$$

We thus obtain the following solitary-wave solution iteratively upon choice of $r_1 = 1$, $r_2 = 1$, $\alpha = 1$ and $\beta = 0.0001$ as follows:

$$u_1(x, t) = 2.449489 \times 10^{-8} ct \operatorname{Sech}(0.01x) \operatorname{Tanh}(0.01x) \\ \times (-2499.99 + 1. \operatorname{Sech}(0.01x)^2 + 0.9999996 \operatorname{Tanh}(0.01x)^2),$$

$$u_2(x, t) = t \operatorname{Sech}(0.01x) \left(6.123724 \times 10^{-14} c^2 t \operatorname{Sech}(0.01x)^6 \right. \\ \left. - 0.0000061 \operatorname{Tanh}(0.01x) + (2.44948 \times 10^{-8} - 2.44948) \right] \\ \times 10^{-8} c \operatorname{Tanh}(0.01x)^3 - 3.061862 \times 10^{-11} c^2 t \operatorname{Tanh}(0.01x)^4 \\ + 1.2247 \times 10^{-14} c^2 t \operatorname{Tanh}(0.01x)^6 \\ + c^2 t \operatorname{Sech}(0.01x)^4 (-1.530931 \times 10^{-10} - 1.592168 \\ \times 10^{-13} \operatorname{Tanh}(0.01x)^2) + \operatorname{Sech}(0.01x)^2 \operatorname{Tanh}(0.01x) (2.44948 \\ \times 10^{-8} + 1.2247448 \times 10^{-7} c + 5.5113519 \\ \times 10^{-10} c^2 t \operatorname{Tanh}(0.01x) - 2.08206 \times 10^{-13} c^2 t \operatorname{Tanh}(0.01x)^3) \right)$$

and so on. The other calculated terms give the closed form solution as

$$u(x, t) = \alpha + \gamma \operatorname{Sech}[k(x + at)]. \quad (23)$$

We illustrate the plots of the residual errors E_m against c with the optimal value of c approaches 0.556 in Fig. 7.

The results obtained by the proposed scheme are in good conformity with that of the standard ADM, listed in Table 4. In addition, the profiles of the solitary wave at $t = 0.5$ and 1.0 are compared in Figs. 8 and 9.

Example 4. Here, the coupled mKdV equation is considered as

$$u_t = 3(uv)_x - 3u^2 u_x + \frac{1}{2} u_{xxx} + \frac{3}{2} v_{xx} - 3\lambda u_x,$$

$$v_t = -3vv_x - 3u^2 v_x - u_{xxx} - 3v_x + 3\lambda v_x, \quad (24)$$

with initial data

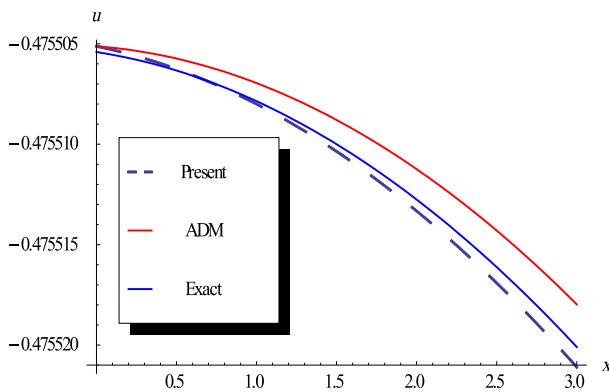


Fig. 8. The plots of the exact, proposed method and the ADM solutions, respectively at $t = 0.5$.

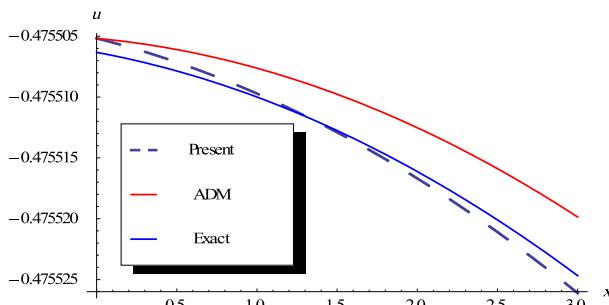


Fig. 9. The plots of the exact, proposed method and the ADM solutions, respectively at $t = 1.0$.

$$u(x, 0) = \frac{b_1}{2k} + k \operatorname{Tanh}(kx),$$

$$v(x, 0) = \frac{\lambda}{2} \left(1 + \frac{k}{b_1} \right) + b_1 \operatorname{Tanh}(kx). \quad (25)$$

Eq. (24) becomes a generalized KdV equation for $u = 0$ and the mKdV equation for $v = 0$, respectively. At this point, we may choose $b_1 = 1; k = 1; \lambda = 1$. Therefore using the recursive relation in Eq. (10) and the Adomian's polynomials in Eq. (11) gives the solution components as follows:

$$u_0(x, t) = \frac{1}{2} + \operatorname{Tanh}(x),$$

$$v_0(x, t) = 1 + \operatorname{Tanh}(x),$$

$$u_1(x, t) = -\frac{1}{4} ct \operatorname{Sech}(x)^2,$$

$$v_1(x, t) = -\frac{1}{4} ct \operatorname{Sech}(x)^2,$$

$$u_2(x, t) = -\frac{1}{16} c^2 t^2 \operatorname{Sech}(x)^2 \operatorname{Tanh}(x),$$

$$v_2(x, t) = -\frac{1}{16} c^2 t^2 \operatorname{Sech}(x)^2 \operatorname{Tanh}(x),$$

$$u_3(x, t) = \frac{1}{2} + \operatorname{Tanh}(x) - \frac{1}{16} ct \operatorname{Sech}(x)^2 (4 + ct \operatorname{Tanh}(x)),$$

$$v_3(x, t) = 1 + \operatorname{Tanh}(x) - \frac{1}{16} ct \operatorname{Sech}(x)^2 (4 + ct \operatorname{Tanh}(x)),$$

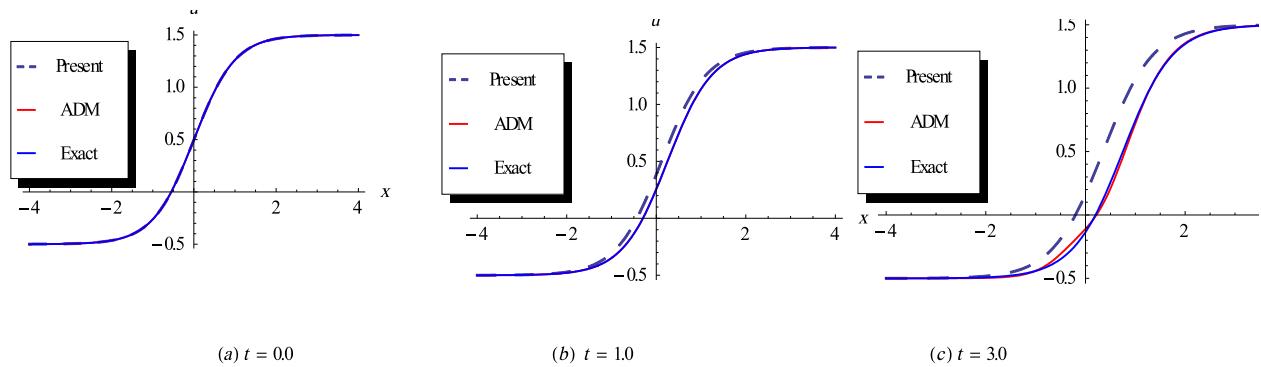


Fig. 10. The plots of the exact, proposed method and the ADM solution, respectively at $t = 0.0, 1.0$ and 3.0 .

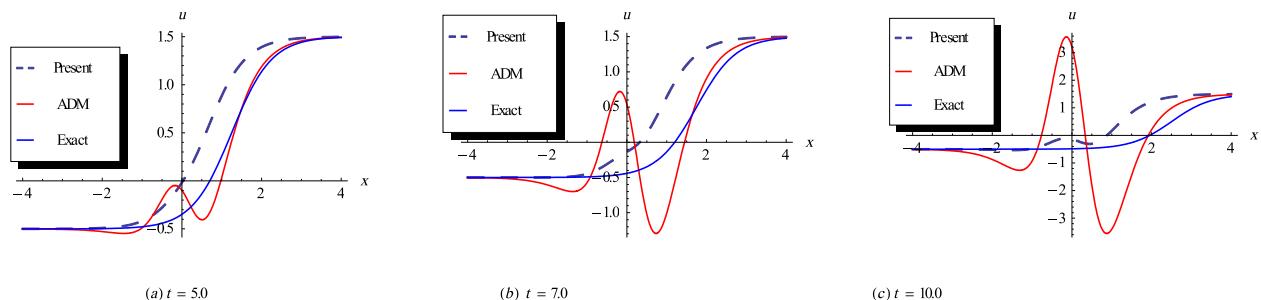


Fig. 11. The plots of the exact, proposed method and the ADM solution, respectively at $t = 5.0, 7.0$ and 10.0 .

and so on. The behavior of the two solutions given by the presented scheme and that of the ADM for certain values of time is depicted and compared with the exact solutions in Figs. 10 and 11. Besides, the exact solutions of $u(x, t)$ and $v(x, t)$ were previously given in (Schiesser, 1994) as

$$u(x, t) = \frac{1}{2} + \operatorname{Tanh}(x + at), \quad v(x, t) = 1 + \operatorname{Tanh}(x + at),$$

$$a = \frac{1}{4} \left(-4k^2 - 6\lambda + \frac{6k\lambda}{b_1} + \frac{3b_1^2}{k^2} \right)$$

5. Conclusion

In conclusion, we have proposed a method based on the recent modification of Adomian decomposition method (ADM) by Zhang and Liang to numerically solve several evolution equations with nonlinearities; particularly the Korteweg de Vries equations. This modification depends on embedding a convergence parameter in the Adomian components. The obtained approximate solutions have been compared to the solutions of the classical AD with the aid of the Mathematica software. The investigated examples show that better accuracy was achieved by the proposed method when compared with those by the ADM for the same number of solution components. The proposed scheme should be considered for extension to other classes of partial differential equations arising in engineering applications.

References

- Abdel-Gawad, H.I., Tantawy, M., Inc, M., Yusuf, A., 2018. On multi-fusion solitons induced by inelastic collision for quasi-periodic propagation with nonlinear refractive index and stability analysis. *Mod. Phys. Lett. B* 1850353.
- Aliyu, T., Shaikh, A.A., Qureshi, S., 2018. Development of a nonlinear hybrid numerical method. *Advan. Differential Equa. Control Processes* 19, 257–285.
- Ansari, M.Y., Shaikh, A.A., Qureshi, S., 2018. Error bounds for a numerical scheme with reduced slope evaluations. *J. Appl. Enviro. Biolog. Sci.* 8, 67–76.
- Adomian, G., 1983. Stochastic Systems. Academic Press, New York.
- Adomian, G., 1986. A new approach to the heat equation—an application of the decomposition method. *J. Math. Anal. Appl.* 113, 202–209.
- Adomian, G., 1988. A review of the decomposition method in applied mathematics. *J. Math. Anal. Appl.* 135, 501–544.
- Adomian, G., 1991. Solving frontier problems modelled by nonlinear partial differential equations. *Comput. Math. Appl.* 22, 91–94.
- Adomian, G., Rach, R., 1992. Noise terms in decomposition series solution. *Comput. Math. Appl.* 24, 61–64.
- Aly, E.H., Ebaid, A., Rach, R., 2012. Advances in the Adomian decomposition method for solving two-point nonlinear boundary value problems with Neumann boundary conditions. *Comp. Math. Appl.* 63, 1056–1065.
- AlQarni, A.A., Banaja, M.A., Bakodah, H.O., Mirzazadeh, M., Biswas, A., 2016. Optical solitons with coupled nonlinear Schrodinger's equation in birefringent nanofibers by Adomian decomposition method. *J. Comput. Theor. Nanoscience* 13, 5493–5498.
- Bakodah, H.O., 2012. A comparison study between a Chebyshev collocation method and the Adomian decomposition method for solving linear system of Fredholm integral equations of the second kind. *JKAU Sci.* 24, 49–59.
- Bakodah, H.O., 2012. Some modifications of Adomian decomposition method applied to nonlinear system of Fredholm integral equations of the second kind. *Int. J. Contemp. Math. Sci.* 7, 929–942.
- Bakodah, H.O., Al-Zaid, N.A., Mirzazadeh, M., Zhou, Q., 2017. Decomposition method for solving Burgers' Equation with Dirichlet and Neumann boundary conditions. *Optik* 130, 1339–1346.
- Bakodah, H.O., Banaja, M.A., AlQarni, A.A., Alshaery, A.A., Younis, M., Zhou, Q., Biswas, A., 2015. Optical solitons in birefringent fibers with Adomian decomposition method. *J. Comput. Theor. Nanoscience* 12, 5846–5853.
- Bakodah, H.O., Darwish, M.A., 2013. Solving Hammerstein type integral equation by new discrete Adomian decomposition methods. *Math. Problems Eng.*, 1–5
- Bakodah, H.O., Mufti, R.S., Biswas, A., 2016. Error estimates of nonlinear algebraic equations by modified Adomian decomposition method. *J. Comput. Theor. Nanoscience* 13, 1–6.
- Banaja, M.A., AlQarni, A.A., Bakodah, H.O., Zhou, Q., Moshokoa, S.P., Biswas, A., 2017. The investigation of optical solitons in cascaded system by improved Adomian decomposition scheme. *Optik* 130, 1107–1114.
- Bulut, H., Belgacem, F.B.M., Baskonus, H.M., 2013. Partial fractional differential equation systems solutions by Adomian decomposition method implementation. In: 4th International Conference on Mathematical and Computational Applications (ICMCA), Manisa, Turkey.
- Ebaid, A., 2011. A new analytical and numerical treatment for singular two-point boundary value problems via the Adomian decomposition method. *J. Comput. Appl. Math.* 235, 1914–1924.
- Ebaid, A., Aljoufi, M.D., Wazwaz, A.M., 2015. An advanced study on the solution of nanofluid flow problems via Adomian's method. *Appl. Math. Lett.* 46, 117–122.
- Duan, J.S., Rach, R., 2011. A new modification of the Adomian decomposition method for solving boundary value problems for higher order nonlinear differential equations. *Appl. Math. Comput.* 218, 4090–4118.
- Duan, J.S., Rach, R., Wang, Z., 2012. On the effective region of convergence of the decomposition series solution. *J. Algorithms Comput. Technol.* 7, 227–247.
- Duan, J.S., 2010. An efficient algorithm for the multivariable Adomian polynomials. *Appl. Math. Comput.* 217, 2456–2467.
- Inc, M., Yusuf, A., Aliyu, A.I., Baleanu, D., 2018. Investigation of the logarithmic-KdV equation involving Mittag-Leffler type kernel with Atangana-Baleanu derivative. *Physica A* 506, 520–531.
- Liao, S., 2010. An optimal homotopy-analysis approach for strongly nonlinear differential equations. *Commun. Nonlinear. Sci. Numer. Simulat.* 15, 2003–2016.
- Nuruddeen, R.I., Muhammad, L., Nass, A.M., Sulaiman, T.A., 2018. A review of the integral transforms-based decomposition methods and their applications in solving nonlinear PDEs. *Palestine J. Math.* 7, 262–280.
- Qureshi, S., Ramos, H., 2018. L-stable explicit nonlinear method with constant and variable step-size formulation for solving initial value problems. *Int. J. Nonlinear Sci. Numer. Simula.*
- Rach, R., 1987. On the Adomian (decomposition) method and comparisons with Picard's method. *J. Math. Anal. Appl.* 128, 480–483.
- Sabi'u, J., Jibril, A., Gadu, A.M., 2018. New exact solution for the (3+1) conformable space-time fractional modified Korteweg-de-Vries equations via Sine-Cosine Method. *J. Taibah Univ. Sci.*
- Schiesser, W.E., 1994. Method of lines solution of the Korteweg-de Vries equation. *Comput. Math. Appl.* 28, 147–154.
- Soliman, A.A., Abdou, M.A., 2008. The decomposition method for solving the coupled modified KdV equations. *Math. Comp. Model.* 47, 1035–1041.
- Syam, M., 2005. Adomain decomposition method for approximating the solution of the Korteweg-de Vries equation. *Appl. Math. Comput.* 162, 65–73.
- Varley, E., Seymour, B.R., 1998. Simple derivation of the N-Soliton solutions to the Korteweg-de Vries equation. *SIAM J. Appl. Math.* 58, 904–911.
- Wazwaz, A.M., 1999. A reliable modification of Adomian decomposition method. *Appl. Math. Comput.* 102, 261–270.
- Wazwaz, A.M., 2001. Blow-up for solutions of some linear wave equations with mixed nonlinear boundary conditions. *Appl. Math. Comput.* 123, 133–140.
- Yusuf, A., Inc, M., Aliyu, A.I., 2018. Fractional solitons for the nonlinear Pochhammer-Chree equation with conformable derivative. *J. Coupled Syst. Multiscale Dyn.* 6, 158–162.
- Yusuf, A., Inc, M., Aliyu, A.I., Baleanu, D., 2018. Solitons and conservation laws for the (2+1)-dimensional Davey-Stewartson equations with conformable derivative. *J. Adv. Phys.* 7, 167–175.
- Zhang, X., Liang, S., 2016. A New Analytic Method with a Convergence-control Parameter for Solving Nonlinear Problems. Cornell University Lib.. arXiv:1609.01550v1 [Math.NA].